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AN APPROACH TO THE NUMERICAL MODELLING
OF CUMULUS-SCALE MOTIONS

by

Richard Arthur Anawalt

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THESIS

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OF CUMULUS-SCALE MOTIONS

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Richard Arthur Anawalt

June 1969

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AN APPROACH TO THE NUMERICAL MODELLING
OF CUMULUS-SCALE MOTIONS

by

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Submitted in partial fulfillment of the
requirements for the degree of
MASTER OF SCIENCE IN METEOROLOGY

from the

NAVAL POSTGRADUATE SCHOOL
June 1969

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ABSTRACT

A numerical model which utilizes the isobaric vorticity equation is developed and applied to cumulus-scale data. The model, together with a modified version of the cumulus convection model of Weinstein and Davis, is applied to data obtained from the National Severe Storms Laboratory in Norman, Oklahoma. The calculations yield real time predictions for height, temperature and relative humidity at seven pressure levels, which are then used as input to the cumulus convection model to obtain vertical profiles of various parameters at specified grid points.

Some results of the calculations are presented along with suggestions for further testing and improvement. The results indicate that further modifications to the approach used are necessary in order to provide more accurate forecasts. Values of the individual terms in the vorticity equation are presented as computed from the observed mesoscale data.

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THEORY

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TABLE OF SYMBOLS AND ABBREVIATIONS

a	Threshold value below which conversion from cloud water to hydrometeor water does not take place (gm m^{-3}).
a'	Threshold value below which conversion from cloud water to hydrometeor water does not take place (gm Kg^{-1}).
C_{pv}	Specific heat of water vapor at constant pressure ($\text{cal gm}^{-1} \text{ }^{\circ}\text{K}^{-1}$).
C_w	Specific heat of liquid water ($\text{cal gm}^{-1} \text{ }^{\circ}\text{K}^{-1}$).
e_s	Saturation vapor pressure in millibars.
f	Coriolis parameter.
f_m	Mean Coriolis parameter for region studied = f_{35} .
g	Acceleration of gravity (m sec^{-1}).
J	Jacobian.
L_f	Latent heat of fusion (cal gm^{-1}).
L_s	Latent heat of sublimation (cal gm^{-1}).
L_v	Latent heat of vaporization (cal gm^{-1}).
m	Kessler's saturation vapor density (gm m^{-3}).
M	Kessler's precipitation content (gm m^{-3}).
p	Pressure in millibars (mb).
Δq_s	Saturation mixing ratio over water minus saturation mixing ratio over ice (gm gm^{-1}).
Q_c	Cloud water content (gm gm^{-1}).
Q_h	Hydrometeor water content (gm gm^{-1}).
R	Gas constant for dry air ≈ 0.28704 ($\text{joules gm}^{-1} \text{ }^{\circ}\text{K}^{-1}$).
R_h	Relative humidity (%).

TABLE OF SYMBOLS AND ABBREVIATIONS, CONTINUED

t	Time.
T	Temperature.
u	Zonal component of horizontal velocity (m sec^{-1}).
v	Meridional component of horizontal velocity (m sec^{-1}).
V	Terminal velocity of hydrometeors (m sec^{-1}).
W	Horizontal wind vector.
W_{χ}	Divergent component of the horizontal wind vector.
W_{ψ}	Rotational component of the horizontal wind vector.
w	Vertical component of velocity in (x, y, z, t) coordinate system (m sec^{-1}).
z	Height of isobaric surface in meters.
e	Ratio of gas constants $\approx .622$.
ζ	Vertical component of relative vorticity (sec^{-1}).
ρ	Air density (gm m^{-3}).
Φ	Geopotential = gz.
ψ	Stream function ($\text{m}^2 \text{sec}^{-1}$).
ω	Vertical component of velocity in (x, y, p, t) coordinate system = $\frac{dp}{dt}$ (mb hr $^{-1}$ or mb sec $^{-1}$).
∇	"Del" operator = $i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y}$.
∇^2	Laplacian operator = $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$.
∇^2	Finite difference form of the Laplacian operator.
\bar{X}	Average value of any parameter X between two pressure levels.
(\bar{X})	Average value of any parameter X with respect to temperature.

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1. INTRODUCTION

The National Severe Storms Laboratory (NSSL) located at Norman, Oklahoma, has been gathering vast amounts of data during the thunderstorm season for the past several years. Many studies have been conducted with the aid of this data (for example see Barnes, 1968, Newton and Fankhauser, 1964 and Hammond, 1967), however, after an extensive literature search, it appears that no research has been done, or at least reported, involving a numerical modelling approach whereby the distributions of various parameters were analyzed and their values numerically forecast for a future time. Such an approach, if successful, could give distinct advantages in predicting, for example, where hail might be expected, or where the most favorable conditions for tornadoes might be located. In addition, during severe convective activity radiosonde data often terminates near mid-tropospheric levels. Consequently, a model which could reasonably predict the sounding above termination level, when given a complete sounding from the previous observation time, would aid greatly in later research efforts.

Figure 1 shows the region used for this study with a height-field analysis and the reported winds at the nine radiosonde stations in the NSSL upper-air network (see Appendix A for a description of the NSSL upper-air network). Assuming the data to be free of errors, the non-geostrophic nature of the flow on this scale is immediately evident, and the figure emphasizes the need for intensive research into the dynamic mechanisms involved under such circumstances. In an effort to evaluate the possibilities of numerical weather prediction for such small-scale data, this paper uses a numerical prediction model based on the complete isobaric vorticity equation.

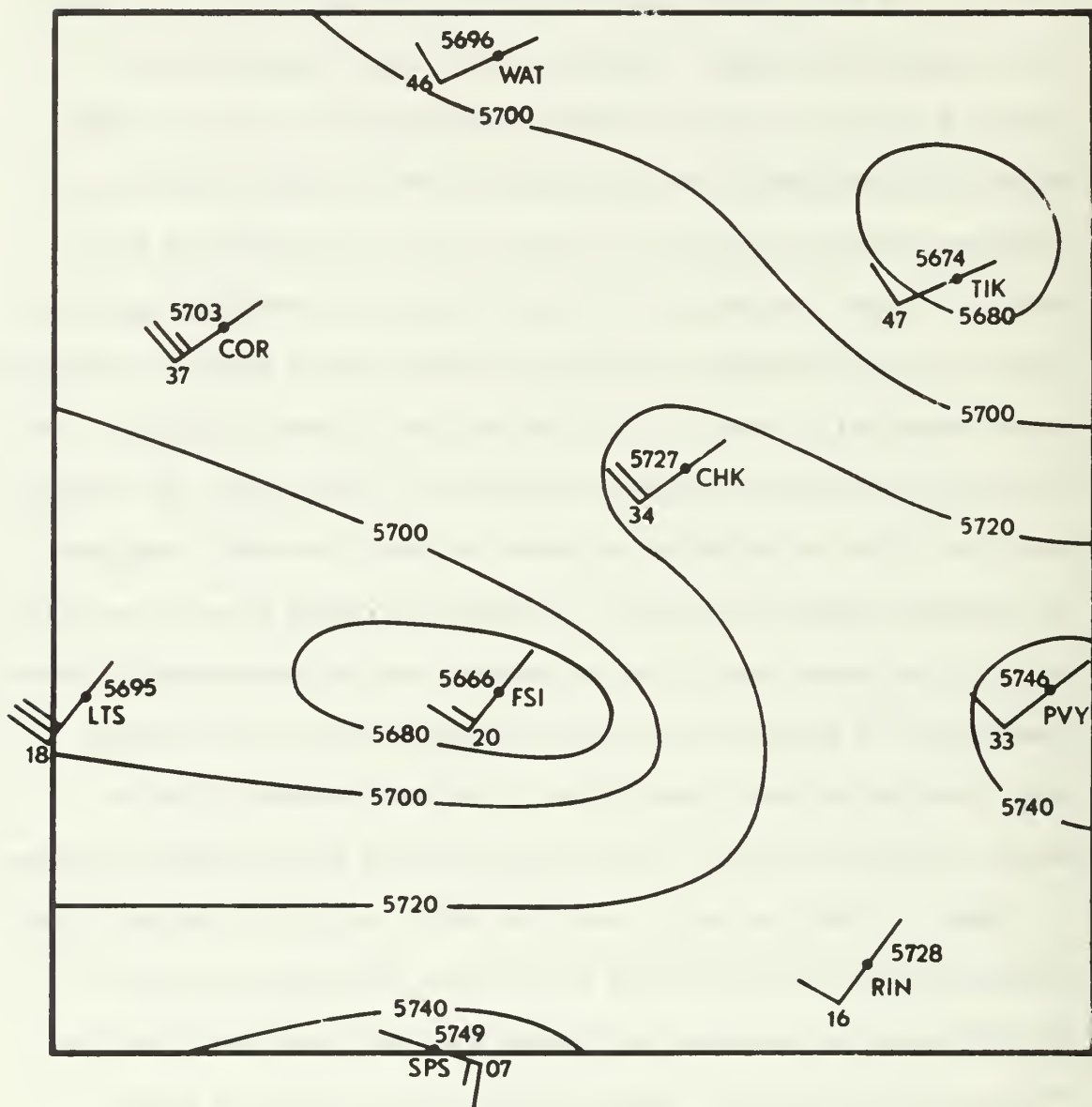


Figure 1: Height analysis with wind directions and speeds (m sec^{-1}) at nine radiosonde stations of NSSL network for the 500-mb level at 1700Z on 30 May 1967 which illustrates the non-geostrophic nature of the data.

The resulting forecasts are then used as input to a cumulus convection model developed by Weinstein and Davis¹ (1967) to provide an aid to severe storm forecasting.

The parameterized cumulus convection model by Weinstein and Davis (1967) calculates cloud top, profiles of vertical velocity, temperature excess (that is, the temperature of a parcel of air in the cloud minus the temperature of a parcel of air in the environment), hydro-meteor water content and updraft radius of the cloud as well as rainfall amount and duration. The principles of the above calculations are based on the entrainment principles of Stommel (1947) and the results of Kessler (1967). With a standard radiosonde sounding at mandatory and significant levels beginning at cloud base, the W.D. model transforms the sounding to one which is spaced in equal increments of height. After the user specifies certain boundary conditions such as ice nucleation temperature, radius of the base of the cloud and conversion and collection rates (Kessler, 1967), the model specifies cloud extent. One use of this model is to compare the effect of seeding a given cloud or allowing it to develop naturally. This is accomplished by varying the ice nucleation temperature.

The numerical model presented in this paper predicts heights, temperatures and relative humidities at individual grid points for seven isobaric levels (950 mb, 850 mb, 700 mb, 500 mb, 300 mb, 200 mb, and 150 mb) so that predicted information can be used as input to a modified subroutine form of the W.D. model to yield the vertical

¹The model developed by Weinstein and Davis has been referred to as the Pennsylvania State model. In this paper, it will be abbreviated as the W.D. model.

profiles previously discussed, under the assumption that cloud base is at 850 mb. However, the present version, as described in this paper, does not provide for vertical coupling of the seven isobaric levels after the first time step.

Predictions were made at 1.5-hour intervals utilizing the IBM OS/360 (MVT versions 15, 16) computer. The 1.5-hour prediction interval was selected to correspond to the times when verification data were available from the NSSL radiosonde network.

The grid used for the computations is a 24 x 24 X-Y grid oriented 000-180, 090-270, having a mean latitude of 35N. The grid distance is 5 nautical miles with the grid arranged so that the nine upper-air stations are either located exactly on a grid point or are so close to one that they may be approximated to be at the nearest grid point. The farthest distance away from a grid point for a particular station was 2 nm. They were assumed to be located at the nearest grid point to eliminate the need for an interpolation scheme in the objective analysis used for initialization of the data as described in Section 4.

2. MODIFICATIONS TO THE W.D. MODEL

While working with and studying the W.D. model, it became apparent that several modifications could be made to it that would provide more accuracy. Although most of the modifications were of a minor nature, three of the modifications are considered to be significant improvements over the original model. These three modifications will be presented in detail while the remainder of the changes will be summarized at the end of this section since they are considered to be self-explanatory.

Kessler (1967) proposed the equation

$$\frac{dM}{dt} = - \frac{dm}{dt} = k_1 (m-a) \quad (2-1a)$$

to model the conversion of cloud water to hydrometeor water. Here, k_1 is an empirical constant which is equal to $.001 \text{ sec}^{-1}$ when $m > a$ and is zero for $m < a$. The quantity "a" is defined as a threshold value and refers to the cloud water content. If the cloud water content is greater than this threshold value, then conversion of cloud water to hydrometeor water will occur. If cloud water amounts less than the threshold value are present, then conversion of cloud water to hydrometeor water will not take place. This equation is incorporated into the W.D. model (Weinstein and Davis, 1967) in the form

$$\frac{dQ_h}{dt} = - \frac{dQ_c}{dt} = k_1 (Q_c - a') \quad (2-1b)$$

where k_1 is defined in a similar manner as it was in (2-1a). Kessler (1967) indicated that measurements have been performed which indicate that cloud water amounts greater than 1 gm m^{-3} are usually associated

with the production of precipitation. Furthermore, he arbitrarily suggested the value of 0.5 gm m^{-3} for the threshold value. Weinstein and Davis (1967) chose $a' = 0.5 \text{ gm Kg}^{-1}$ for their threshold value.

It is proposed to consider this threshold value to be a function of density. This approach thereby associates a different threshold value with each level by making it a function of height. Using this approach, the conversion of cloud water content to precipitation in the W.D. model has been modified to

$$\frac{dQ_h}{dt} = k_1(Q_c - a/\rho) \quad (2-2)$$

where ρ = air density. By defining k_1 as it was in (2-1a), conversion of cloud water to precipitation is modelled to occur when $(Q_c - \frac{a}{\rho}) > 0$.

Computations have been made to determine the variation of the quantity a/ρ . These values show a maximum of approximately $.00137 \text{ gm gm}^{-1}$ in the upper levels and a minimum value of $.00076 \text{ gm gm}^{-1}$ near cloud base.

The second significant modification is the inclusion of a linear variation of the latent heats of vaporization (L_v) and fusion (L_f) with temperature. Since

$$L_s = L_v + L_f \quad (2-3)$$

and L_s may be considered constant while L_f is a function of temperature, then L_v is also a function of temperature. The values of L_f (List, 1951) are 48.6 cal gm^{-1} at -50°C and 79.7 cal gm^{-1} at 0°C . The latent heat of vaporization varies from $629.3 \text{ cal gm}^{-1}$ at -50°C to $597.3 \text{ cal gm}^{-1}$ at 0°C . Since L_s varies by only 1 cal gm^{-1} from 0°C to -40°C , its value was taken as constant and equal to 678 cal gm^{-1} . The

variation of only one of the latent heats on the right side of equation (2-3) needs to be computed since the value of the second one can be obtained by subtraction.

In order to approximate the variation of the latent heat of vaporization as a linear function of temperature, it was necessary to determine an optimum value for the difference between the specific heats of water and ice. Integrating the formula

$$\frac{dL_v}{dT} = C_{pv} - C_w \quad (2-4)$$

and applying the mean value theorem for integrals yields

$$L_v = L_{v_o} + (\overline{C_{pv} - C_w}) (T - T_o). \quad (2-5)$$

Utilizing the values of L_v (List, 1951) for various temperatures (T) and with $T_o = -40C$, the value of $(\overline{C_{pv} - C_w})$ that satisfies equation (2-5) was computed for each temperature at 10 degree intervals from -40C to +40C. A plot of $(\overline{C_{pv} - C_w})$ versus temperature is shown in Figure 2. Based upon the results it was decided to select the value of -.62340 as an "optimum" value for $(\overline{C_{pv} - C_w})$ and consider it to be a constant. This value is considered to be a representative value between -40C and 0C, and corresponds exactly to the value at -10C.

The final significant modification involves the values used for the constants when computing saturation vapor pressure (Weinstein and Davis, 1967). Integrating the Clausius-Clapeyron equation one obtains the following for the saturation vapor pressure over water:

$$\begin{aligned} \ln e_s = & - \frac{\epsilon [L_{vo} - T_o (\overline{C_{pv} - C_w})]}{RT} + \frac{\epsilon (\overline{C_{pv} - C_w}) \ln T}{R} \\ & + \frac{\epsilon [L_{vo} - T_o (\overline{C_{pv} - C_w})]}{RT_o} - \frac{\epsilon (\overline{C_{pv} - C_w}) \ln T_o}{R} + \ln e_{so}. \end{aligned} \quad (2-6)$$

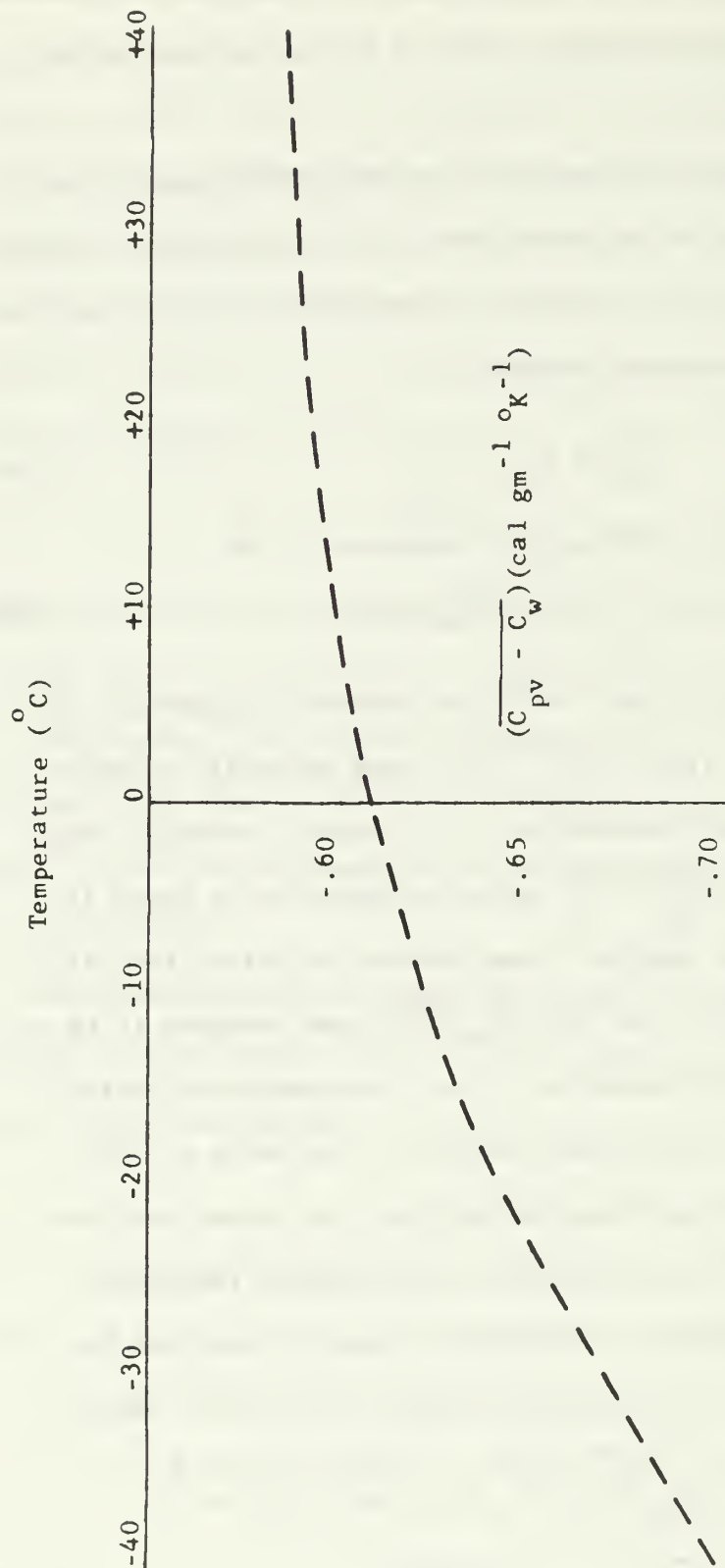


Figure 2: Variation of $(C_{pv} - C_w)$ with temperature.

Simplifying yields

$$\ln e_s = -\frac{A'}{T} - B' \ln T + C'$$

where

$$A' = \frac{\epsilon [L_{vo} - T_o (\overline{C_{pv}} - \overline{C_w})]}{R} = 6958.9262$$

$$B' = -\frac{\epsilon (\overline{C_{pv}} - \overline{C_w})}{R} = 5.65567$$

and

$$C' = \frac{\epsilon}{R} \left[\frac{L_{vo}}{T_o} - (\overline{C_{pv}} - \overline{C_w}) (1 + \ln T_o) \right] + \ln e_{so} = 59.01383$$

when $T_o = 233.16K$, $(\overline{C_{pv}} - \overline{C_w}) = -0.62340 \text{ cal gm}^{-1}K^{-1}$

and $L_{vo} = 621.7 \text{ cal gm}^{-1}$.

Similarly, when freezing occurs one obtains

$$\ln e_s = \frac{-A''}{T} + C'' \quad (2-7)$$

where

$$A'' = \frac{\epsilon L_s}{R} = 6151.0205$$

and $C'' = \frac{A''}{T_o} + \ln e_{so} = 24.3277$ when $T_o = 233.16K$ and $L_s = 678 \text{ cal gm}^{-1}$.

The above equations yield the saturation vapor pressure in millibars.

Obviously the number of significant digits shown above exceeds the accuracy of some of the assumptions made to arrive at the values shown.

Rounding to five significant figures could be done with no loss of accuracy, however they have been retained here as they were computed on the IBM OS/360 computer.

There were several minor changes made to the W.D. model. These included changing the value being used for the acceleration of gravity from 9.87 m sec^{-2} to a more exact form obtained by computing the value from the formula (List, 1951) expressing g as a function of latitude, altering the values used for the constants for conversion from degrees

Celsius to degrees Kelvin, conversion from meters to feet, and the value of the gas constant for dry air. In addition, the section of the program dealing with corrections due to the wind shear were completely eliminated which consequently allowed a subroutine (subroutine CHEQ) to be eliminated. Other minor changes included removing factors of 10 so that all pressure values are carried in units of millibars, introducing symbols wherever possible to simplify future modification and testing, changing the estimate of the updraft velocity at cloud base from 2.0 to 0.5 m sec^{-1} , and modifying certain format statements to improve the headings on the output. Finally, many soundings were run through the final modified version of the W.D. model to ensure that all changes had been properly programmed and to ensure that the results from the model were reasonable.

3. MATHEMATICAL DEVELOPMENT OF A FORECAST MODEL

This section describes the model used for this research. Although it is not consistent in many respects, it did provide a means by which forecasts of height at 1.5-hour intervals could be obtained for input to the W.D. model as well as provide a method by which the values of the various terms in the prediction equation could be ascertained. A description at the end of this section describes an attempt to set up the prediction equation in a consistent manner after the original model was running properly, however, the more consistent method failed to provide reasonable forecasts.

A. THE VORTICITY EQUATION

Following the procedures given by Haltiner and Martin (1957), the isobaric vorticity equation may be written as

$$\frac{d}{dt} (\zeta + f) = - (\zeta + f) \nabla \cdot \mathbf{W} + \left(\frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p} - \frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p} \right). \quad (3-1a)$$

Since the entire grid used for this research covers only 2 degrees of latitude, the Coriolis parameter can be considered constant, taking on the value it would have at the mean latitude of the grid. Using this approximation and expanding the left side of (3-1a) yields

$$\frac{\partial \zeta}{\partial t} = - \left(\omega \frac{\partial \zeta}{\partial p} + \mathbf{W} \cdot \nabla \zeta + (\zeta + f_m) \nabla \cdot \mathbf{W} + \frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p} - \frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p} \right). \quad (3-1b)$$

In order to use the W.D. model at 1.5-hour intervals, height of the base of the clouds is needed from which computations can be started. To obtain these heights, (3-1b) will be used as a height forecasting equation analogous to methods applied to large scale motions in the atmosphere by applying the balance equation. An alternative method without using the balance equation for initialization was tried and is discussed at the end of this section.

Since the application of (3-1b) to height forecasting will be done to obtain information regarding the orders of magnitude of the various terms and to obtain heights which can be used as cloud base information for input to the W.D. model, no attempt was made to set up a consistent set of equations. To obtain an equation which could be used for height forecasting and evaluation of the various terms with mesoscale data, the relative vorticity on the left side of (3-1b) and in the divergence term was replaced by $\nabla^2 \psi$. In addition, only the rotational wind components were used for the horizontal advection term, that is,

$$u = - \frac{\partial \psi}{\partial y} \text{ and } v = \frac{\partial \psi}{\partial x}. \quad (3-2)$$

This was done because the present version has no provision for predicting the divergent component of the wind (W_χ). Future modifications should provide the means by which W_χ can be used in addition to forecasts of the rotational component (W_ψ). The vorticity in the vertical advection term as well as divergence and the u and v components of the wind in the twisting term were computed based on the observed winds at the initial time and held constant for 1.5 hours. Keeping values constant in the $\frac{\partial}{\partial p}$ terms greatly simplifies programming considerations since 3-minute time steps were used and to update these values at each time step would involve working with each level for only one time step before proceeding to the next level. Unfortunately, holding ζ in the vertical advection term and the u and v components of the wind in the twisting term constant for 1.5 hours causes the seven levels to lack vertical coupling. Future modifications should attempt to correct this deficiency. After making these substitutions and assumptions, the resulting prediction equation becomes

$$\frac{\partial}{\partial t}(\nabla^2 \psi) = - \left(\omega \frac{\partial \zeta}{\partial p} - \frac{\partial \psi}{\partial y} \frac{\partial \zeta}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \zeta}{\partial y} + (\nabla^2 \psi + f_m) \nabla \cdot W + \frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p} - \frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p} \right). \quad (3-3)$$

Assuming continuity in ψ , the term on the left side of (3-3) may be written as $\nabla^2 \frac{\partial \psi}{\partial t}$, and the equation then becomes a Poisson-type equation which can be solved by relaxation procedures as described by Haltiner and Martin (1957). The forecasting procedure, being one of successive iterations, will thus allow all terms involving ψ in (3-3) to vary at each time step.

To obtain a stream field for (3-3), consider the balance equation (Charney, 1955):

$$\nabla^2 \psi + \frac{1}{f} (\nabla \psi \cdot \nabla f) - \frac{2}{f} \left[\left(\frac{\partial^2 \psi}{\partial x \partial y} \right)^2 - \frac{\partial^2 \psi}{\partial x^2} \frac{\partial^2 \psi}{\partial y^2} \right] = \frac{\nabla^2 \Phi}{f}. \quad (3-4)$$

Since the Coriolis parameter is to be considered constant over the grid, the second term of (3-4) drops out.

As discussed by Charney (1955) and Bolin (1955), equation (3-4) will maintain its elliptic character only if

$$\frac{\nabla^2 \Phi}{f} + \frac{f}{2} > 0, \quad (3-5)$$

that is, only if the geostrophic relative vorticity is greater than $-f/2$. As shown in Appendix B, heights would have to be read at least to the nearest centimeter on this size grid in order to satisfy the requirements of (3-5) and still allow negative geostrophic relative vorticities. Consequently, (3-4) was forced to be elliptic by writing the non-linear terms in the form

$$\left(\frac{\partial^2 \psi}{\partial x \partial y} \right)^2 = - \frac{\partial u}{\partial x} \frac{\partial v}{\partial y}, \quad (3-6a)$$

and

$$\frac{\partial^2 \psi}{\partial x^2} \frac{\partial^2 \psi}{\partial y^2} = - \frac{\partial v}{\partial x} \frac{\partial u}{\partial y}. \quad (3-6b)$$

Making this substitution into (3-4) and taking into account that the Coriolis parameter is to be considered a constant, (3-4) becomes

$$\nabla^2 \psi = \frac{g}{f_m} \nabla^2 z - \frac{2}{f_m} J(u, v), \quad (3-7)$$

where J represents the Jacobian operator, and u and v are the components of the observed winds. Equation (3-7) is then an elliptic equation as it stands and does not require that the geostrophic relative vorticity be greater than $-f/2$. The ψ field can be initialized by the approximation $\psi = \frac{gz}{f_m}$, and then equation (3-7) can be relaxed to convergence before using (3-3) to predict the ψ fields at a later time. Similarly, once the ψ fields have been forecast ahead, the height fields can be recovered by making the approximation $z = \frac{f_m \psi}{g}$ and then utilizing (3-7) in the relaxation procedure until convergence is attained.

Convergence was defined as occurring when (3-7) balanced within 0.5 meters. When converting from height fields to stream fields this requires an epsilon of approximately $58612 \text{ m}^2 \text{ sec}^{-1}$, whereas epsilon becomes 0.5 m when converting from stream fields to height fields. These values arise from the approximation $\psi = (g/f_m)z$. In retrospect, the above convergence criteria may be too large for this scale of motion. In the future, more stringent criteria (such as 0.1 m) should be invoked.

As mentioned at the beginning of the section, attempts were made to correct the deficiencies of the prediction equation. This involved utilizing the exact relation $\nabla^2 \psi = \zeta$ as the method of obtaining the initial ψ field instead of using the balance equation and to substitute $\nabla^2 \psi$ for ζ in the horizontal advection term. In addition, the convergence criteria was changed to $10000 \text{ m}^2 \text{ sec}^{-1}$, in units of ψ , when converting from ζ to ψ .

The balance equation was used to recover height fields from the forecast fields of ψ with convergence criteria of 0.1 m. Two methods of initializing the ψ field were tried; they were $\psi = \frac{\overline{gz}}{f_m}$ and $\psi = \frac{gz}{f_m}$,

where \bar{z} represents the standard atmosphere value of the height of the respective isobaric surfaces. When using \bar{z} for the initialization, computations were carried out without difficulty, however the resulting forecast height fields were very flat as might have been anticipated. However, the forecast fields were in error by as much as 300 m at some levels. When using z for initialization, instability resulted in the horizontal advection term which caused an exponential growth of this term to values too large for the computer to handle after only 12-15 time steps (36-45 minutes ahead in time). The $\nabla^2 \psi$ part of the horizontal advection term is believed to be the cause of the instability but this was not investigated in detail. However, one attempt to investigate the instability was made by shortening the time step to one minute. This delayed the instability somewhat which would indicate the problems were caused by computational instability.

B. VERTICAL MOTION COMPUTATIONS

In order to solve (3-3), vertical motion values are required for the vertical advection of relative vorticity and for the twisting term. Vertical motions for each of the seven levels were computed by means of the continuity equation in the form

$$\frac{\partial \omega}{\partial p} = - \overline{\nabla \cdot \mathbf{V}} \quad (3-8)$$

where the bar symbol denotes a vertical average through the layer. This method of computation utilizes the measured wind fields and does not allow the ω values to be updated for forecast times since the present model has no provisions for obtaining the divergent component of the wind. The vertical boundary conditions imposed were

$$\omega_{1000} = 0.0 \quad (3-9a)$$

and $\omega_{150} = 0.0 \quad (3-9b)$

which also imply that $\nabla \cdot W_{1000} = - \sum_{p=950}^{150} \nabla \cdot W.$

Vertical motion computations were made by working downward from 150 mb through 300 mb and upward from 1000 mb through 700 mb. Vertical motion at 500 mb was then computed as the average of the values at 300 mb and 700 mb. It is recognized that using a backward or forward finite difference form for $\frac{\partial w}{\partial p}$ could at times lead to computational difficulties, however the method used here does allow the vertical motion at a given level to be affected by not only the divergence at that level and the level above, but also takes into consideration the divergence field at the level below.

Computations for the various levels were as follows (symbols such as W150 and W950 were chosen to conform to the symbols used in the computer program (subroutine VERTMO) shown in Appendix C):

(1) 150 mb

$$W150=0.0$$

(2) 200 mb ($\Delta p > 0$)

$$W175=W150-\Delta p \overline{\nabla \cdot W}_{200-150}$$

$$W225=W150-\Delta p \overline{\nabla \cdot W}_{300-150}$$

$$W200=(W175+W225)/2$$

(3) 300 mb ($\Delta p > 0$)

$$W250=W200-\Delta p \overline{\nabla \cdot W}_{300-200}$$

$$W350=W200-\Delta p \overline{\nabla \cdot W}_{500-200}$$

$$W300=(W250+W350)/2$$

(4) 1000 mb

$$W1000=0.0$$

(5) 950 mb ($\Delta p < 0$)

$$W_{925} = W_{1000} - \Delta p \overline{\nabla \cdot W}_{850-1000}$$

$$W_{975} = W_{1000} - \Delta p \overline{\nabla \cdot W}_{950-1000}$$

$$W_{950} = (W_{925} + W_{975}) / 2$$

(6) 850 mb ($\Delta p < 0$)

$$W_{825} = W_{950} - \Delta p \overline{\nabla \cdot W}_{700-950}$$

$$W_{900} = W_{950} - \Delta p \overline{\nabla \cdot W}_{850-950}$$

$$W_{850} = (2(W_{825}) + W_{900}) / 3$$

(7) 700 mb ($\Delta p < 0$)

$$W_{675} = W_{850} - \Delta p \overline{\nabla \cdot W}_{500-850}$$

$$W_{725} = W_{850} - \Delta p \overline{\nabla \cdot W}_{500-950}$$

$$W_{700} = (W_{675} + W_{725}) / 2$$

(8) 500 mb

$$W_{500} = (W_{300} + W_{700}) / 2$$

The values obtained for vertical motion by this method are discussed in Section 5.

C. TEMPERATURE AND RELATIVE HUMIDITY FORECAST SCHEME

In order to utilize the W.D. model as an aid in severe storm forecasting, forecasted values of temperature and relative humidity are necessary in sounding form as input to the model.

Although a model designed for operational use should definitely be internally consistent, such as the vertical consistency of temperature fields with the geopotential fields, this study involved determining some of the problems that would develop when applying numerical forecasting techniques to mesoscale data. Consequently, the scheme for obtaining temperature forecasts which follows is not dynamically consistent with what has preceded, but it does provide a method of

obtaining a temperature-field forecast which could be used as input to the W.D. model. As discussed in Section 6, future studies should improve on this method.

Four different methods of forecasting temperature were tested and the results were then compared to verifying data. The method which verified best was selected as the forecast method for temperature, which was simply

$$\frac{\partial T}{\partial t} = - 1/2 W \cdot \nabla T, \quad (3-10)$$

where the values for the wind components were taken as the measured winds at observation time and were not altered during the forecast interval.

Similar testing of four prediction methods was also done for relative humidity. The results led to

$$\frac{\partial R_h}{\partial t} = - 1/2 W \cdot \nabla R_h \quad (3-11)$$

as the prediction equation for relative humidity where the wind components were again taken as the measured winds at observation time and were held constant for the forecast interval.

Temperature and relative humidity prediction calculations by equations (3-10) and (3-11) were performed in subroutine PROG, a listing of which may be found in Appendix C. Results of the computations, when compared to verifying data, are found in Section 5.

D. BOUNDARY CONDITIONS

To initialize the ψ field, the approximation $\psi \approx g\bar{z}/f_m$ was used as discussed in Section 3A. When the relaxation was performed, the outside row or column of the grid could not be modified due to finite

differencing. As a result, the boundary values of ψ were held at the value obtained by the initialization. The prognostic equation (3-3) was utilized from $x = 3, 22$ and $y = 3, 22$. After initializing the $\frac{\partial \psi}{\partial t}$ field to zero, it was then relaxed for the first time step. This relaxation was performed from $x = 3, 22$ and $y = 3, 22$ as mentioned above and the outer two rows or columns were maintained at their original values of zero. By this means, actual forecast values of ψ were obtained only from $x = 3, 22$ and $y = 3, 22$. After the first time step, the previous relaxed values for $\frac{\partial \psi}{\partial t}$ were used as the initial estimate for the new time step.

After obtaining the 1.5-hour forecast for the ψ field, the conversion back to z fields by means of (3-7) was accomplished by making the initial approximation $z \approx f_m \psi / g$ only for the points where the ψ field had been forecast. This then leaves the outer two rows or columns without forecasted height values. Before relaxing the z field, however, the values at these boundary points were set by performing a linear extrapolation from the fourth row or column outward as follows:

$$\begin{array}{rcl} 4 & . & \\ 3 & . & \\ 2 & . & z_2 = z_3 - (z_4 - z_3) \\ 1 & . & z_1 = z_2 - (z_3 - z_2) \end{array}$$

Once this was accomplished at all of the grid points in the two boundary rows or columns, the z field was relaxed by means of the balance equation from $x = 2, 23$ and $y = 2, 23$.

Horizontal boundary values were not a problem with the vertical motion fields since they were computed from $x = 2, 23$ and $y = 2, 23$.

As mentioned above, the prognostic equation was utilized from $x = 3, 22$ and $y = 3, 22$ so that there were computed values of ω at $x = 2, x = 23, y = 2$ and $y = 23$ to use for the terms $\frac{\partial \omega}{\partial x}$ and $\frac{\partial \omega}{\partial y}$. As stated in (3-9), vertical boundary conditions for ω were that the vertical motion would vanish at 150 mb and 1000 mb, which then implies that the divergence at 1000 mb is equal to the negative of the sum of the divergences from 950 mb to 150 mb.

As will be discussed in Section 4, the initial data fields of height, temperature, relative humidity, wind direction and wind speed were analyzed on the computer by an objective scheme. The analysis method included finite differencing methods such that the outer row or column was not analyzed. These boundary values were set to the same value as the next interior row or column.

The same linear extrapolation technique that was used for the height fields was also used for the boundary conditions of the forecasted values of temperature and relative humidity, but it was only necessary to extrapolate from the second row or column out to the first row or column.

E. TIME STEPS

In the forecasting of the new ψ fields by (3-3), a 3-minute time step was used, thus requiring 30 time steps to obtain a 1.5-hour forecast. The method used was the leap-frog scheme described by Haltiner (upbl.), with a forward time step used for the first step. To eliminate the possibility of obtaining two separate solutions by the leap-frog method, the two solutions were averaged every 10 time steps

(30 minutes) and then one pass through the averaged field was made utilizing a Laplacian smoother with a .04 coefficient for smoothing. The leap-frog forecast method was then continued.

The forecasting of temperature and relative humidity by the simple advection schemes described in Section 3C does not involve any iterative type computations since the wind components were held constant for the forecast interval. Consequently, a single 90-minute forward time step was used. However, future modifications, which should include provisions for obtaining rotational and divergent components of the wind, should also update the wind field with each time step so that forecast values of the wind field could be used for temperature and relative humidity advection.

4. INITIALIZATION OF THE MODEL

In order to commence computations with the previously discussed model, 35 fields of data are needed (7 levels each of height, temperature, relative humidity, wind direction and wind speed). Such a volume of data would be difficult to obtain by manual means for real-time operational use of a model. Therefore, a more realistic approach is to use an objective analysis scheme by the computer.

Although many complicated and very sophisticated analysis schemes can be and have been developed (for example see Hughes, 1967), this was not the purpose of the research reported in this paper. Consequently, a simple objective analysis scheme was devised which uses the Laplacian operator and an eight-point averaging technique in alternating steps. Comparisons of computer analyses and hand analyses showed the patterns to be very similar. The major difference between the two analysis methods was the gradient in the vicinity of closed centers. The hand analyses tended to subjectively spread the gradient out over a much larger area whereas this was not so with the computer analyses. Which analysis is more nearly correct is merely a matter of subjective judgement. With only nine data points within a 14,400 square mile area during conditions of severe storms, this question will probably not be answered until the data network becomes more detailed. Nevertheless, the computer analyses were judged to yield adequate data fields for testing of the computer program.

In order to start the analysis a first guess was needed at each grid point. This was accomplished by dividing the grid into nine "regions of influence" (see Figure 3), each region containing one

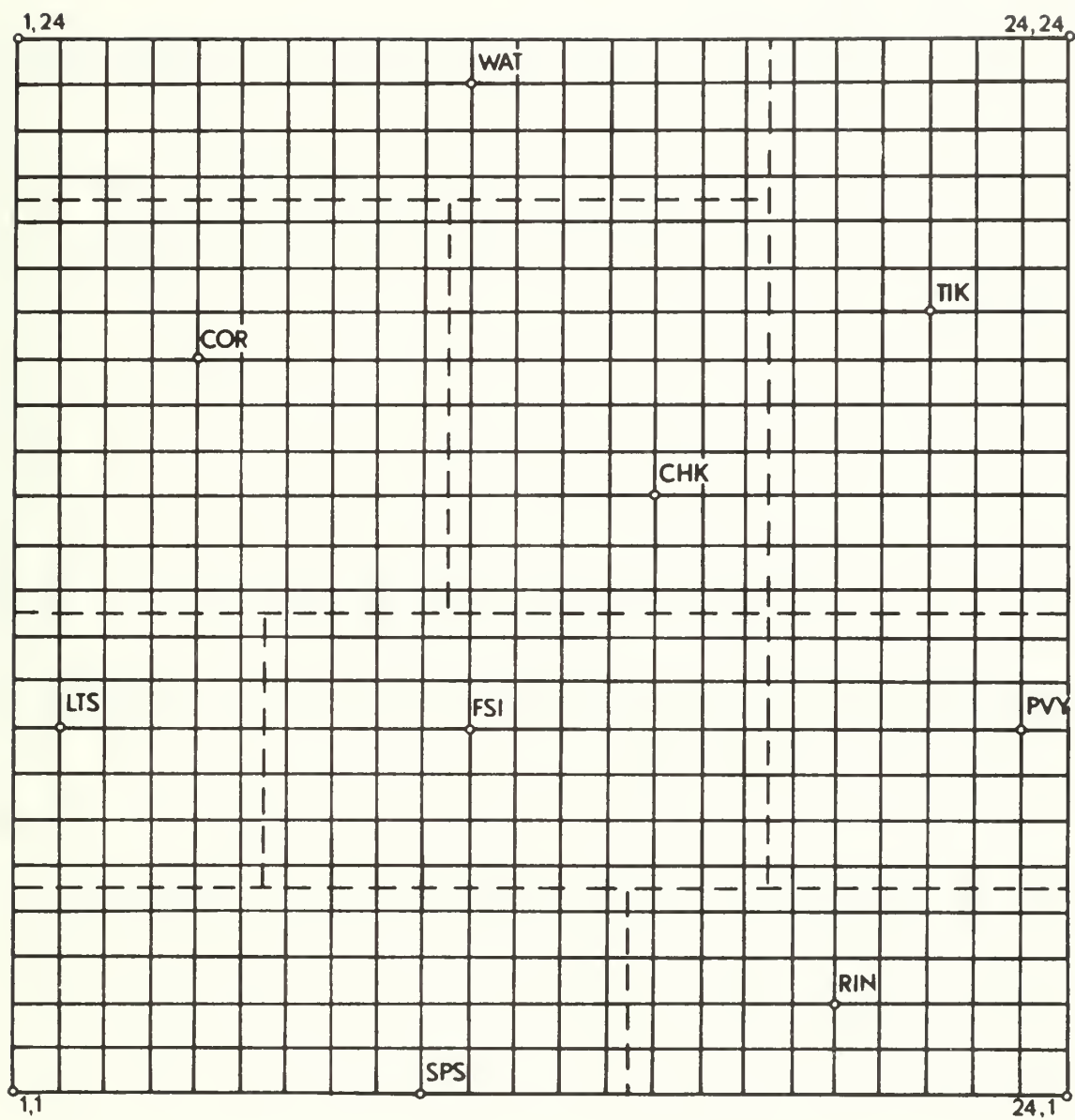
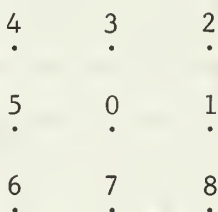


Figure 3: Regions of Influence.

of the nine radiosonde stations. The value of a given parameter at each grid point within a region was then set equal to the known value at the radiosonde station.

With the initial guess completed, the analysis was started. The first step was to alter the value at the grid points (except for the nine data points) by setting the value at that point equal to the average of the surrounding eight points according to the following diagram:



$$P_0 = \frac{P_1 + P_2 + P_3 + P_4 + P_5 + P_6 + P_7 + P_8}{8} \quad (4-1)$$

After making one pass through the grid in this manner, the Laplacian-type scheme was used such that

$$P_0 = P_0 + \delta (P_1 + P_3 + P_5 + P_7 - 4P_0), \quad (4-2)$$

where the value of $\delta = 0.1$ was arbitrarily selected. After making ten passes through the grid in this manner, the eight-point averaging technique was again employed. The entire procedure was repeated in such a manner that the eight-point average was employed 11 times while the Laplacian scheme was utilized 100 times. The resulting analysis, after boundary conditions were set, was the analysis used in the prediction model. The boundary condition imposed was that the value at a grid point on an outside row or column was equal to the value at the adjacent grid point on the next row or column toward the center of the grid.

Figures 4a and 5a show hand analyses for two different levels and Figures 4b and 5b show the corresponding computer objective analyses. A copy of the program for the objective analysis scheme is included in Appendix C.

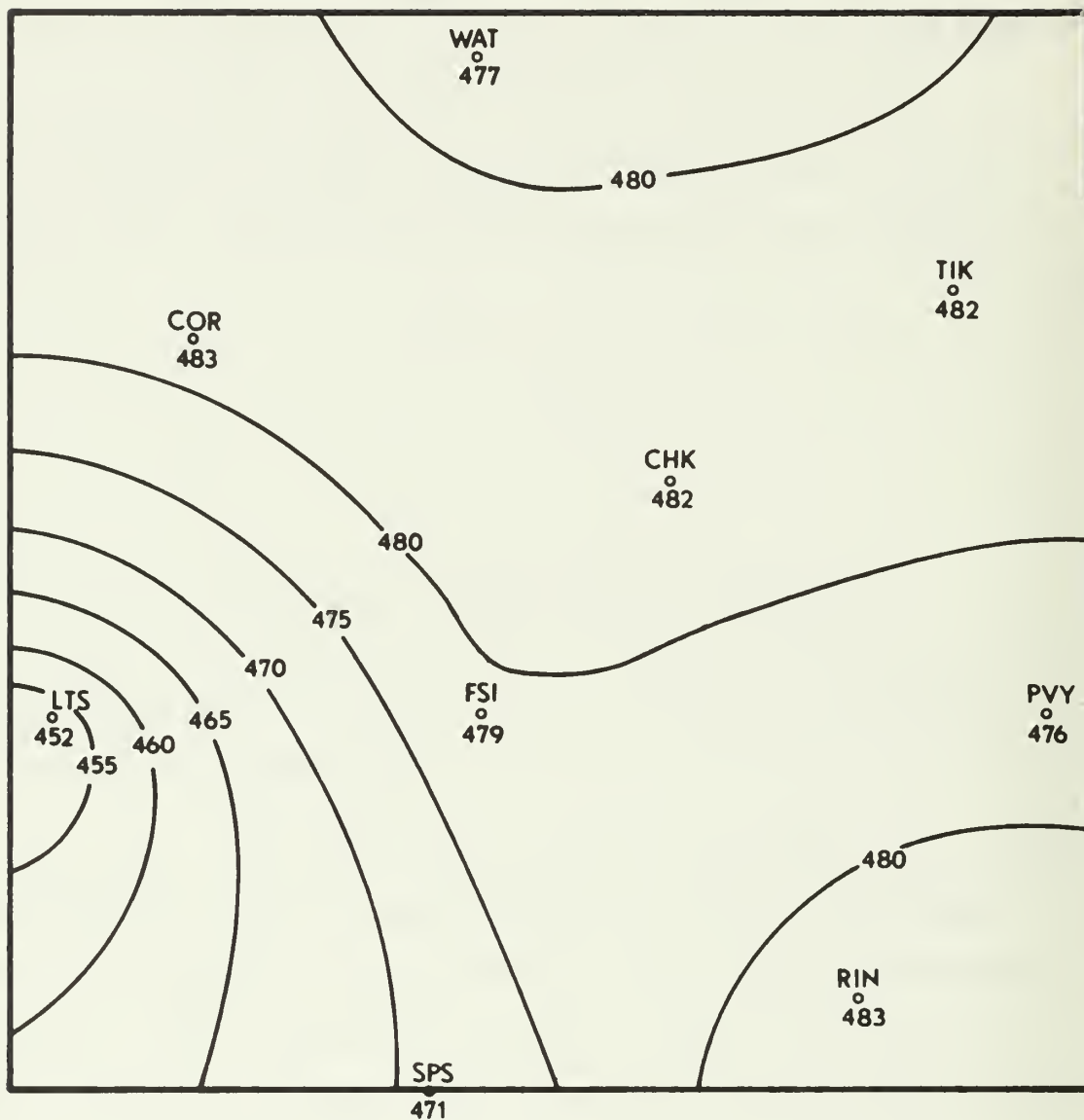


Figure 4a: Hand analysis for 950 mb at 1700Z on 30 May 1967.
Contours every 5 meters.

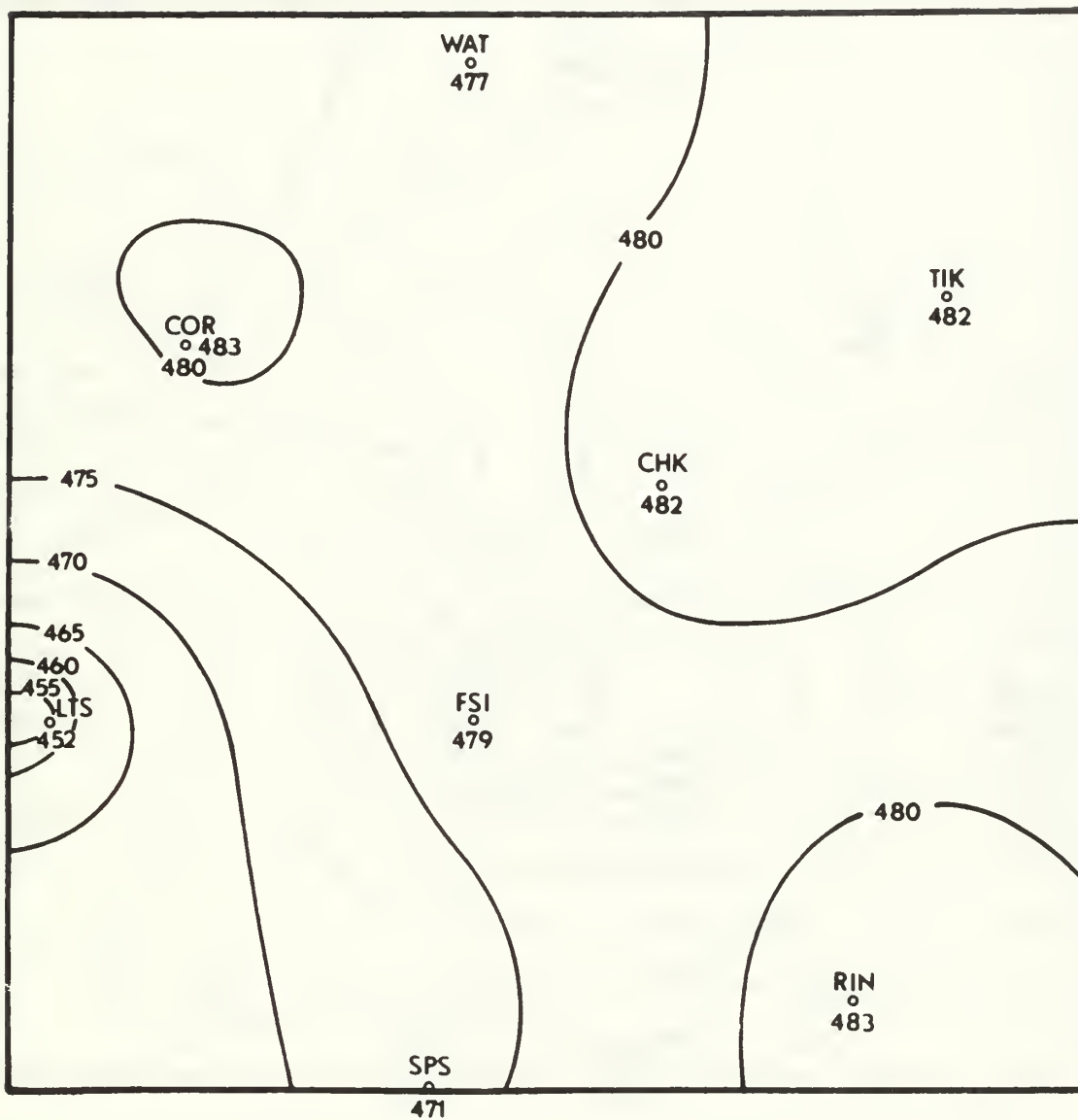


Figure 4b: Computer objective analysis for 950 mb at 1700Z on 30 May 1967. Contours every 5 meters.

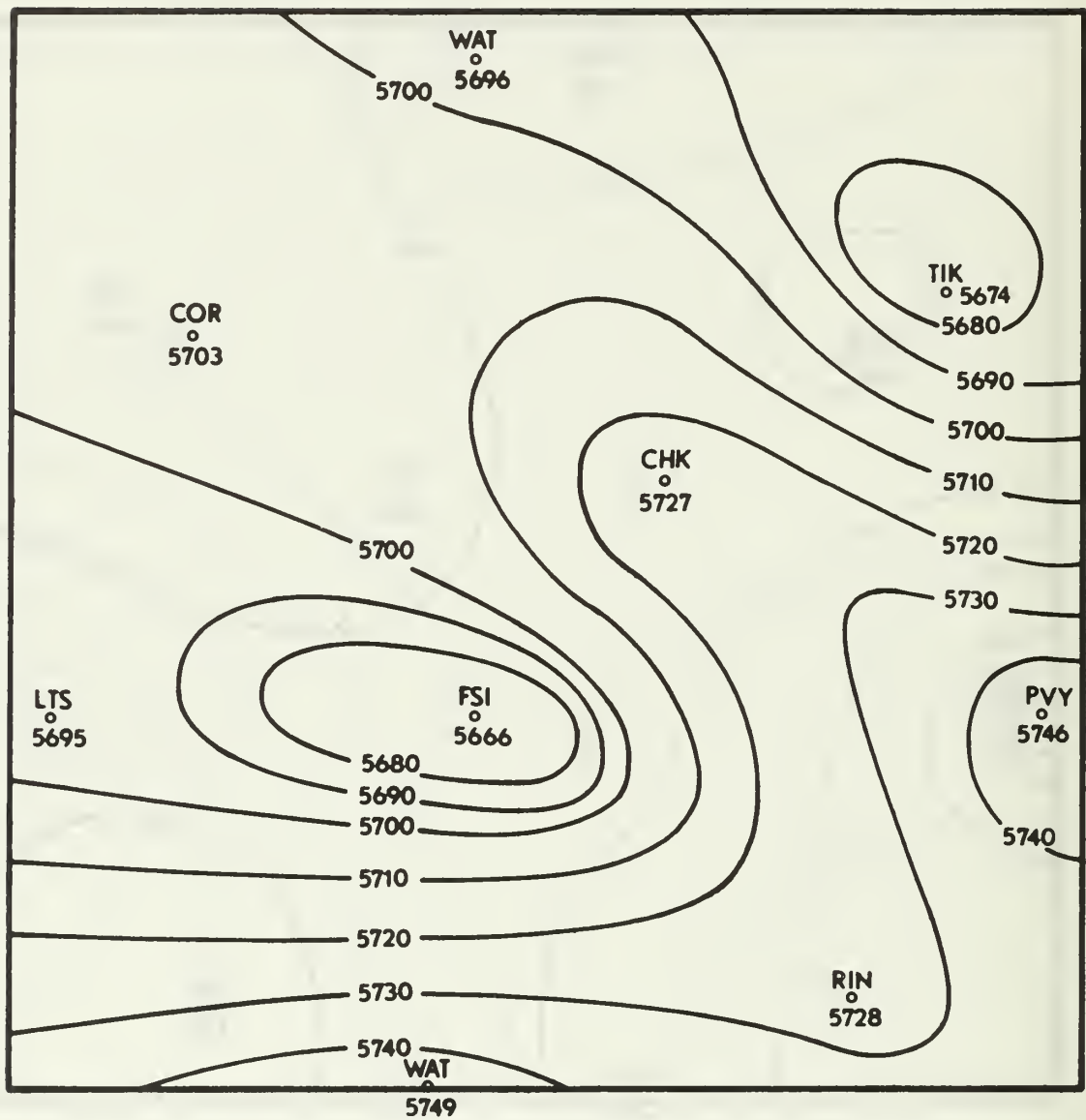


Figure 5a: Hand analysis for 500 mb at 1700Z on 30 May 1967. Contours every 10 meters.

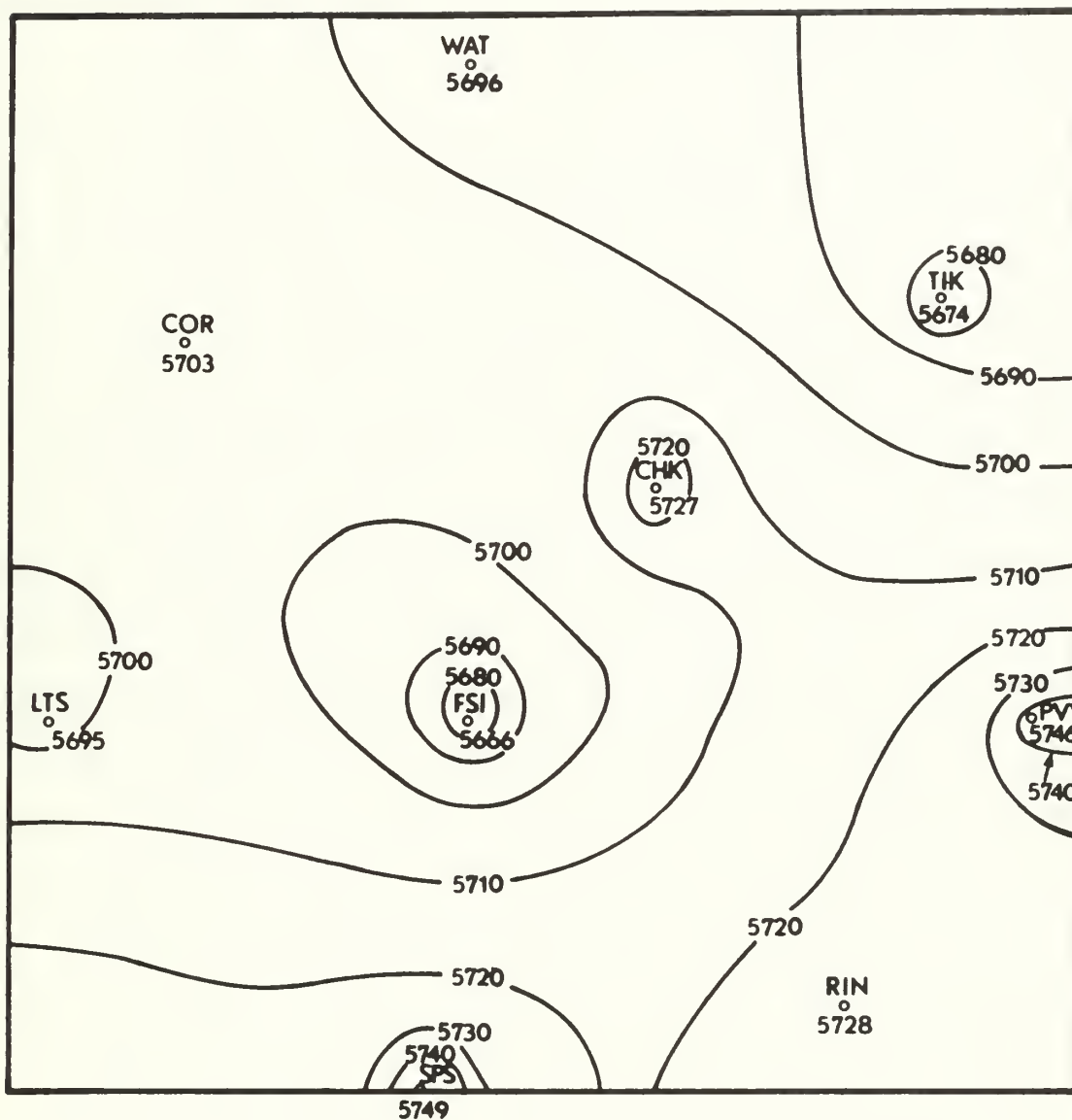


Figure 5b: Computer objective analysis for 500 mb at 1700Z on 30 May 1967. Contours every 10 meters.

5. RESULTS OF THE COMPUTATIONS

Two sets of data have been utilized with the previously described model. The times of the two sets of data are 1700Z, 30 May 1967 and 2130Z, 30 May 1967.

In synoptic-scale models, the terms in equation (3-1a) involving vertical motion and divergence have often been neglected since they may be considered relatively unimportant by scale analysis (Haltiner, upbl.) in the large scale motions of the atmosphere. However, when dealing with cumulus-scale data, this may not be true. In order to evaluate the relative importance of the various terms in the model throughout the entire grid, computations of the orders of magnitude of the various individual terms of equation (3-3) have been performed for both the 2130Z and 1700Z sets of observed data for the seven levels in the vertical described in Section 1. Tables 2 and 3 show the minimum and maximum values of each term based on the observed data, not on the forecast values. Although these tables only show the minimum and maximum values, they do point out the fact that the divergence term is very important. There were many instances when all five terms were of the same order of magnitude, that is, when each term was just as important as each of the other terms. Table 4 shows a sample of the values of each of the five terms as extracted from the computer printouts.

Out of 4800 values tested, $\frac{\partial v}{\partial x} \frac{\partial v}{\partial p}$ and $-\frac{\partial v}{\partial y} \frac{\partial u}{\partial p}$ were of the same sign 2584 times, or about 54% of the time. Therefore, in the cases tested, they were additive more often than not which is further justification that they should be retained for small-scale data.

PRESSURE LEVEL

TERM (UNITS OF 10^{-7})	950		850		700		500		300		200		150	
	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX
$\omega \frac{\partial \zeta}{\partial p}$	-.624	.206	-.258	.270	-1.62	.731	-.259	.286	-.192	.435	-.137	.067	0.0	0.0
$-\frac{\partial \psi}{\partial y} \frac{\partial \zeta}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \zeta}{\partial y}$	-6.04	9.81	-11.4	10.9	-24.3	20.9	-17.9	22.2	-13.8	19.2	-22.2	46.6	-8.16	12.7
$(\nabla^2 \psi + f_m) \nabla \cdot W$	-11.7	11.9	-19.5	14.5	-47.2	55.0	-168.	74.4	-49.4	42.8	-147.	284.	-21.9	29.8
$\frac{\partial w}{\partial x} \frac{\partial v}{\partial p}$	-.369	.701	-.310	.195	-.909	1.41	-.539	.897	-.644	1.21	-.218	.066	0.0	0.0
$-\frac{\partial w}{\partial y} \frac{\partial u}{\partial p}$	-.321	.816	-.492	.351	-.933	1.50	-.246	.491	-1.26	.884	-.158	.163	0.0	0.0

Table 2: Minimum and maximum values of components of equation (3-3) for 1700Z data for all seven isobaric levels.

PRESSURE LEVEL

TERM (UNITS OF 10^{-7})	950		850		700		500		300		200		150	
	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX	MIN	MAX
$\omega \frac{\partial \zeta}{\partial p}$	-.532	.143	-.174	.074	-.203	.120	-.066	.108	-.130	.107	-.031	.068	0.0	0.0
$-\frac{\partial \psi}{\partial y} \frac{\partial \zeta}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \zeta}{\partial y}$	-5.82	8.70	-4.67	8.51	-13.4	10.9	-3.43	16.7	-12.7	22.0	-97.1	49.3	-47.2	24.5
$(\nabla^2 \psi + f_m) \nabla \cdot \mathbf{V}$	-34.3	18.2	-22.0	23.6	-44.2	104.	-17.3	21.0	-47.3	164.	-200.	257.	-226.	68.9
$\frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p}$	-.432	1.06	-.302	.992	-.356	.510	-.273	.347	-.564	1.47	-.160	.082	0.0	0.0
$-\frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p}$	-.318	.503	-.366	.363	-.590	1.13	-.142	.100	-.438	.395	-.126	.105	0.0	0.0

Table 3: Minimum and maximum values of components of equation (3-3) for 2130Z data for all seven isobaric levels.

$\omega \frac{\partial \zeta}{\partial p}$	$-\frac{\partial \psi}{\partial y} \frac{\partial \zeta}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \zeta}{\partial y}$	$(\nabla^2 \psi + f_m) \nabla \cdot W$	$\frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p}$	$-\frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p}$
0.18E-08	-0.15E-08	-0.21E-09	-0.25E-08	0.22E-08
0.20E-08	0.24E-07	0.22E-09	-0.14E-08	0.64E-09
0.19E-08	0.56E-07	0.11E-08	-0.95E-09	-0.33E-09
0.11E-08	0.11E-06	0.32E-08	0.32E-09	0.91E-09
-0.30E-10	0.16E-06	0.88E-08	0.46E-08	0.36E-08
-0.20E-08	0.19E-06	0.14E-07	0.70E-08	0.21E-08
-0.43E-08	0.13E-06	0.20E-07	0.79E-08	-0.31E-09
-0.61E-08	0.34E-07	0.20E-07	0.66E-08	0.95E-10
-0.73E-08	-0.16E-08	0.15E-07	0.51E-08	-0.84E-09
-0.69E-08	-0.11E-07	0.68E-08	0.43E-08	-0.63E-09
-0.75E-08	-0.16E-07	0.24E-08	0.28E-08	-0.62E-10
-0.67E-08	0.99E-08	-0.45E-09	-0.35E-09	0.98E-11
-0.73E-08	0.28E-07	-0.14E-08	-0.19E-08	0.59E-09
-0.40E-08	0.27E-07	-0.19E-08	-0.16E-07	0.29E-08
-0.11E-08	0.23E-07	-0.28E-08	0.43E-08	0.45E-08
0.92E-09	0.11E-07	-0.34E-08	0.61E-08	0.39E-08
0.13E-08	0.50E-08	-0.34E-08	0.31E-08	0.12E-08
0.84E-09	0.18E-08	-0.30E-08	0.21E-09	-0.14E-08
0.30E-09	0.14E-08	-0.20E-08	-0.28E-09	-0.37E-08
0.41E-09	-0.41E-07	-0.12E-08	-0.82E-09	0.79E-08
0.13E-08	-0.10E-07	-0.13E-09	-0.12E-08	0.31E-08
0.10E-08	0.24E-07	0.14E-09	-0.23E-08	0.24E-09
0.56E-09	0.54E-07	0.59E-09	-0.30E-08	-0.22E-09
-0.25E-09	0.11E-06	0.26E-08	-0.26E-08	0.46E-08
-0.15E-08	0.22E-06	0.80E-08	0.26E-08	0.75E-08

Table 4: Sample of values of the components of equation (3-3) at 950 mb for 1700Z data.

Vertical motion computations yielded considerably different values for the 1700Z data and the 2130Z data. Table 5 shows the maximum upward and downward vertical motions for each isobaric level at each of the times. Although the values in Table 5 may, at first, seem somewhat large, a transformation by the formula

$$\omega \approx -\rho \, g w \quad (5-1)$$

where ρ is replaced by p/RT from the equation of state to yield

$$\omega \approx -\frac{p}{RT} \, g w, \quad (5-2)$$

shows that a value of $\omega = 200 \text{ mb hr}^{-1}$ upward corresponds to a value of $w \approx .86 \text{ m sec}^{-1}$ upward (at 500 mb) which is certainly a reasonable value for cumulus convection. Although values this large were only present at a few points, further research is needed to determine what reasonable values are for this scale of motion. In general, the vertical motion fields showed fairly strong upward vertical motions throughout the grid at 1700Z and correlated fairly well with the squall-line that was present at that time. There were a few areas in which the correlation was not too good but this may possibly be due to the fact that the nine radiosondes were not all launched simultaneously and a difference of 15 minutes between launch times could possibly result in the bad correlation regions. Although the radar pictures indicated that dissipation of the squall-line had begun at 2130Z by the fact that isolated echoes were present rather than a large mass of clouds, the vertical motion fields did not show extensive regions of downward motions except at 300 mb. The 2130Z values were, however, considerably less in intensity than they were at 1700Z which might have been anticipated from the results shown in Table 5.

Pressure Level Up/Down/Time	950	850	700	500	300	200
Max. Upward 1700Z	-98	-74	-220	-105	-108	-31
Max. Downward 1700Z	+78	+86	+219	+55	+76	+21
Max. Upward 2130Z	-48	-30	-110	-39	-53	-18
Max. Downward 2130Z	+52	+46	+90	+72	+78	+10

Table 5: Maximum upward and downward vertical motions for each isobaric level at 1700Z and 2130Z on 30 May 1967 in mb hr^{-1} .

The significance of the much larger values of vertical motion at 1700Z as compared to 2130Z was not investigated thoroughly, however larger values of divergence were evident (as would be expected) at 1700Z. This was particularly true at 200 mb, and this would tend to be propagated downward to 300 mb and 500 mb by the computation procedure.

The simple method of forecasting temperature resulted in values which verified slightly better than a persistence forecast. The average absolute error (not considering whether the error was an overestimate or an underestimate) for 97 persistence temperature forecasts was 1.8C while for 105 50% temperature advection forecasts the average absolute error was 1.4C. The reason for the different number of verifications is due to the fact that some of the 1700Z and 2130Z radiosonde data terminated before reaching 150 mb. Verifications were performed only at the nine radiosonde stations, and only at each of the seven discrete isobaric levels.

The simple method of forecasting relative humidity resulted in values which verified essentially the same as a persistence forecast. The average absolute error for 97 persistence relative humidity forecasts was 18.0% while for 105 50% relative humidity advection forecasts the average absolute error was 18.2%.

The results of the height forecasts did not verify as well as was expected. This is undoubtedly due in part to the fact that additional refinements are necessary in the prediction method to obtain a consistent set of equations. Table 6 compares the average absolute errors of persistence and model 1.5-hour forecasts for both the 1700Z and the 2130Z sets of data. As can be seen in the table, the model did perform slightly better than persistence 50% of the time.

AVERAGE ABSOLUTE ERROR (METERS)

TYPE OF FORECAST (VERIFYING TIME)	950	850	700	500	300	200	150
Persistence (1830Z)	4.5	8.4	18.2	26.0	26.8	31.4	34.2
Model (1830Z)	4.4	7.1	8.8	28.5	28.1	31.8	31.0
Persistence (2300Z)	13.1	8.9	10.6	6.5	31.1	44.5	53.1
Model (2300Z)	12.9	9.1	7.2	14.9	37.5	32.0	58.5

Table 6: Average absolute errors of height fields (meters) for persistence and model at 1700Z and 2130Z for all seven isobaric levels.

As previously mentioned, the W.D. model yields profiles of vertical velocity, hydrometeor water content and temperature excess. In addition, total rain expected from the cloud, duration of the precipitation, updraft area, and a profile of updraft radius are included in the results of the computations. Figure 6 shows a sample of the graphical output of the W.D. model. These results are printed out in graphical form by the two subroutines (GRPHCL and GRAPH) which were included with the W.D. model (Weinstein and Davis, 1967). The horizontal axis (shown in Figure 6 at 1/2 scale) shows the scale to be used for the various parameters and the vertical axis on the left shows the corresponding heights in meters, with the bottom figure (in this case 1642) being the first increment (200 meters) above cloud base. In the actual computer output of Figure 6, the lines are represented by letters, such that the vertical velocity curve is a series of the letter "W", hydrometeor water content of the letter "Q", temperature excess of the letter "T", updraft radius of the letter "R", and the zero axis for the temperature excess curve is represented by the symbol "\$".

The W.D. model has a provision for setting a constant value for the updraft radius or for allowing it to vary within a cloud. However, for this research the updraft radius was always considered to be constant and equal to 1 km. The profile shown in Figure 6 was computed based upon the 1830Z, 30 May 1967 forecasted values of temperature, relative humidity and height of the cloud base, and by assuming that cloud base was at 850 mb for station CHK.

Many soundings were used as input to the W.D. model, and in many cases the values of the temperature excess determined by the model seemed to be too large. Values as high as 4 or 5 degrees Celsius

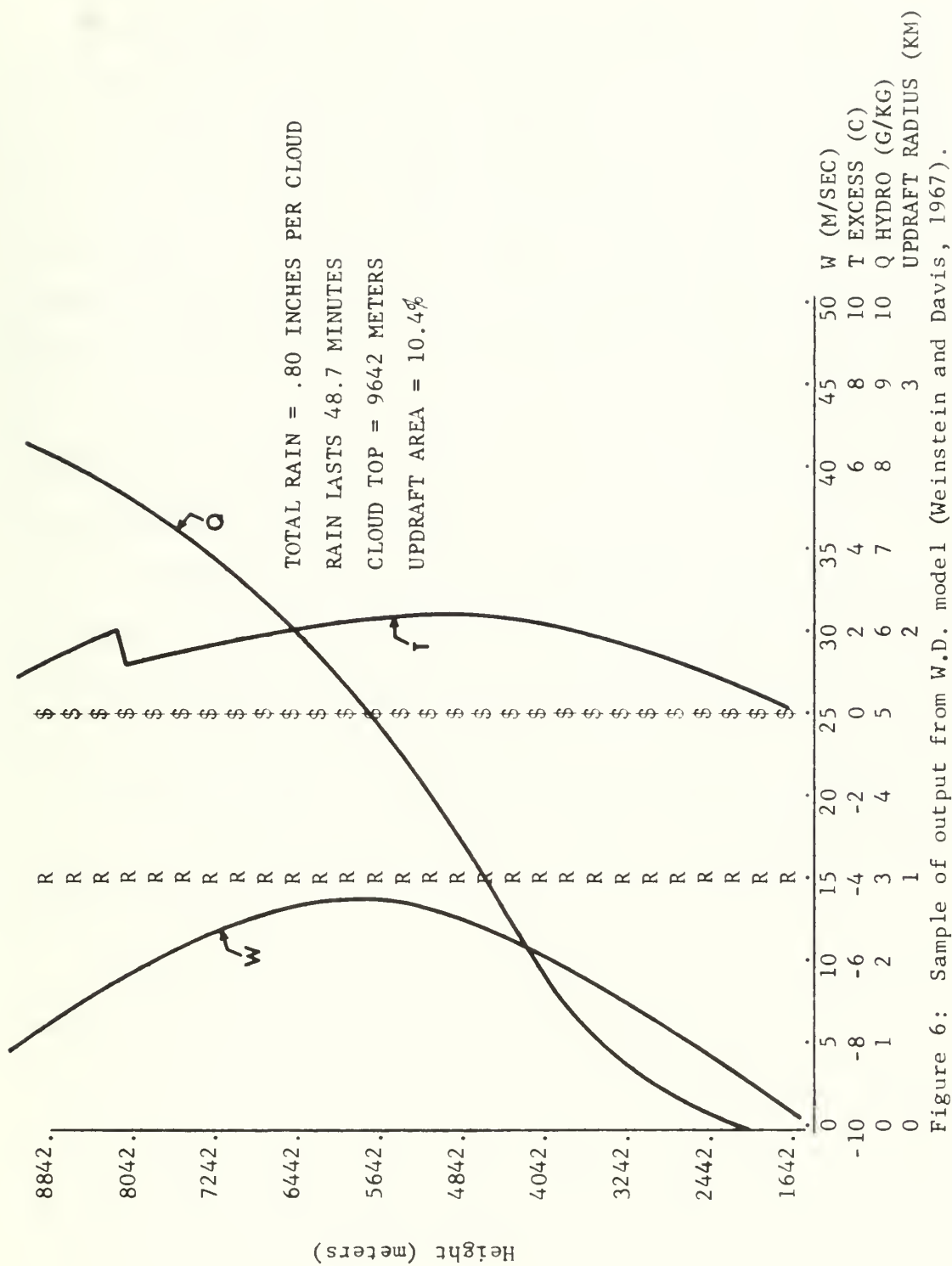


Figure 6: Sample of output from W.D. model (Weinstein and Davis, 1967).

occurred quite often, with an occasional value of 7 or 8 degrees Celsius. This intuitively seems to be too high but additional observations are needed to determine if such values actually exist within cumulus clouds. In addition, apparent discontinuities in temperature excess, as evidenced near 8000 meters on Figure 6, are quite common in the graphical output of the model. These large temperature excess values and temperature excess discontinuities occurred not only with forecast soundings but also with actual observed soundings from the NSSL data.

Figure 6 is based upon an ice nucleation temperature of -25C which is assumed to be representative of natural conditions in the atmosphere. The W.D. model allows for changing the ice nucleation temperature to a value of, say, -6C to represent the conditions that might occur if cloud seeding were to be undertaken. However, this is presently not included in the subroutine version of the W.D. model (subroutine CLOUDS) since the reason for using it as part of a prediction model was to aid in severe storm forecasting.

6. SUGGESTIONS FOR FUTURE DEVELOPMENTS AND IMPROVEMENTS

The first modification to the previously discussed model, which became evident during testing of the model, would be to expand the horizontal grid dimensions to at least 32×32 . Figure 3 illustrates that four of the nine grid points lie within the first two rows of the boundaries. This more or less negates application of the prognostic equations to these grid points due to finite differencing. It is estimated that this improvement would increase the storage requirements of the program by about 50%, and since the program is already quite large, modifications would have to be made to allow for the sharing of some storage locations wherever possible.

At present the model lacks the proper provisions for predicting the complete wind field and for updating it after each time step. A consistent scheme for updating the wind field should be included such that the predicted winds could be recovered every time step. In addition, this would allow predictions further ahead in time as well as allow for updating the vertical motion fields with new divergence fields.

As discussed in Section 3A, the present version of the forecast model does not provide for vertical coupling of the seven isobaric levels since the vorticity in the vertical advection term and the u and v components of the wind in the twisting term were held constant for the 1.5-hour forecast interval. To correct this deficiency would involve performing one time step at a given level before proceeding to the next level. Also, additional storage would be required to save the field which $\frac{\partial}{\partial p}$ operates on for one time step back so that vertical derivatives could be taken at each level on fields which all correspond

to the same time. As far as programming considerations are concerned, this essentially involves nesting the vertical level DO LOOP inside of the time step DO LOOP.

When the model has been modified to allow for vertical coupling of the seven isobaric levels, additional levels in the vertical could be included so that vertical layer depths would not exceed 100 mb. This would require perhaps 10 or 11 levels which would increase the storage requirements of the computer considerably, however, this would allow a much finer detail for the forecasted soundings. In addition, it should aid in the delineation of the vertical velocities at each level.

A consistent dynamical method should be included for the prediction of temperature and relative humidity. The present method is more empirical than dynamical and better verifications should be obtained if a more realistic approach is devised. As an alternative, a better constant (vice 50%) might be tested, particularly in the case of relative humidity where the errors were negative much more often than they were positive.

In subroutine SETUP (see Appendix C), cloud base is assumed to be at 850 mb for entry into subroutine CLOUDS (see Appendix C). A routine could be developed to be included in SETUP such that the lifting condensation level, mixing condensation level or convective condensation level would be computed as the cloud base before entering the W.D. model (subroutine CLOUDS). Comparison of results obtained for cloud base at 950 mb and 850 mb indicate that the cloud base and the associated temperature and relative humidity are very critical to the results obtained from the W.D. model.

Once a complete internally consistent model is set up, calculations out past 1.5 hours could be tried to determine at what point such a model loses its skill. A few runs were made out to 4.5 hours and the results appeared to indicate values of height which were too large, and the temperature fields appeared to become unstable.

Kessler (1967) proposed the formula for the fall speed of precipitation to be of the form

$$V = 5.0939 M^{.125} . \quad (6-1)$$

The corresponding formula in the W.D. model is

$$V = 15.39 Q_h^{.125} . \quad (6-2)$$

The units of M are gm m^{-3} whereas the units of Q_h are gm gm^{-1} . In order to convert M to the units of Q_h , it is necessary to divide by the density of air, ρ . Weinstein and Davis (1967) have apparently selected a representative value for ρ . It is recommended that future modifications of the W.D. model include the density in the computation. Since $M = \rho Q_h$, and $\rho = p/RT$ from the equation of state, it would be a simple change in the programming to use the following formula for the hydrometeor terminal velocity:

$$V = 5.0939 \left(\frac{pQ_h}{RT} \right)^{.125} . \quad (6-3)$$

Since the duration calculation is based on the hydrometeor terminal velocity at cloud top, (6-3) will permit values of V which not only depend on the hydrometeor liquid water content but also will depend on the pressure and temperature at the cloud top.

Finally, it has been suggested to the author by Dr. J. D. Mahlman (personal communication) that the only solution to overcome the many difficulties with this scale of motion would be a direct integration

of the primitive equations. It is believed that no such research has been attempted for mesoscale data and this approach definitely warrants consideration for future attempts at numerical modelling of such scales of motion.

7. SUMMARY

A numerical weather prediction model has been described for use with cumulus-scale data. Together with the W.D. model, it provides forecasts of height, temperature, and relative humidity at seven isobaric levels as well as vertical profiles of vertical velocity, temperature excess and hydrometeor liquid water content at 1.5-hour intervals. The model was tested on two sets of data and yielded forecasts which were, in general, statistically about the same as a persistence forecast. This can probably be attributed in part to the fact that the present version has no provisions for coupling the seven isobaric levels nor for updating the vertical motion fields after each time step. Another factor is that the temperature and relative humidity forecast schemes were more empirical than dynamical, and were not consistent with the vorticity equation. This may account for the lack of skill in forecasting these parameters. Suggestions were made to correct some of these inconsistencies in the event that further research is conducted with the model. However, the possibility of applying numerical forecasting techniques to small-scale data does appear to have possibilities but extensive research is needed in this area before useful forecasts will be obtained.

The model requires approximately 5.5 minutes on the IBM OS/360 (MVT versions 15, 16) computer to yield a 1.5 hour forecast. An additional 14 seconds per field is required for the objective analysis scheme described in Section 4, thus requiring about 13.5 minutes for the analysis of 35 fields of data and the prediction of the aforementioned parameters for 1.5 hours.

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APPENDIX A

THE NSSL RADIOSONDE NETWORK

There were 10 radiosonde stations within the NSSL region considered for this study which are listed below. The grid-point numbers associated with each station are shown on Figure 3, where the stations are shown by their call letters.

Call letters	Name	Location	
		Lat.	Long.
SPS	Sheppard AFB, Texas	33.97	98.48
RIN	Ringling, Oklahoma	34.17	97.58
LTS	Altus, Oklahoma	34.70	99.33
FSI	Fort Sill, Oklahoma	34.65	98.40
PVY	Pauls Valley, Oklahoma	34.70	97.22
COR	Cordell, Oklahoma	35.30	98.97
CHK	Chickasha, Oklahoma	35.10	97.97
OKC	Oklahoma City, Oklahoma	35.40	97.60
TIK	Tinker AFB, Oklahoma	35.42	97.38
WAT	Watonga, Oklahoma	35.85	98.42

All of these stations, except OKC, were used in this study. The reason for not including data from OKC is that their data corresponds to synoptic data times rather than at 1.5-hour intervals.

APPENDIX B

ON THE ELLIPTICITY REQUIREMENT OF THE BALANCE EQUATION

The ellipticity requirement of the balance equation (Charney, 1955 and Bolin, 1955) was given by equation (3-5) which is

$$\frac{\nabla^2 \phi}{f} + \frac{f}{2} > 0. \quad (3-5)$$

Bolin (1955) stated that the data at grid points which did not satisfy (3-5) was modified so that (3-5) was satisfied. This section investigates the extent to which grid distance influences the requirement of (3-5).

Expanding (3-5) results in

$$\frac{\frac{2}{m} g}{d^2 f_m} \nabla^2 z + \frac{f_m}{2} > 0 \quad (B-1)$$

where

∇^2 = finite difference form of the Laplacian operator,

m = map factor ≈ 1.19 ,

g = gravity ≈ 9.80 ,

d = grid distance,

z = height,

and f_m = mean Coriolis parameter = $.836 \times 10^{-4}$.

Upon substitution of the above values for each factor, (B-1) reduces to

$$\nabla^2 z + .00087 d^2 > 0 \quad (B-2)$$

where z is in meters and d is in nautical miles. Substituting a value of $d = 5$ nm yields

$$\nabla^2 z > -.02175. \quad (B-3)$$

By measuring heights to the nearest meter, the only way to satisfy

(B-3) is to have $\nabla^2 z \geq 0$. To satisfy (B-3) on a grid with $d = 5$ nm

requires z to be measured at least to the nearest .01 m, and this would still not guarantee meeting the requirement. It is easily shown with equation (B-2) that the smallest grid distance on which (3-5) can be satisfied, without eliminating the possibility of $\nabla^2 z < 0$ and still measure z to the nearest meter, is 35 nm.

The result from above is what prompted the substitution of the measured wind field into the balance equation in order to force it to be elliptic. Tests were made to determine with what frequency the balance equation ellipticity requirement was satisfied with this scale of data. The results show that out of 6776 points, only 62.4% satisfied the requirement. At 1700Z the requirement was satisfied 59.4% of the time while at 2130Z the figure was 65.5%.

APPENDIX C

This appendix includes a listing of the computer program used for the research described by the preceding sections. In order to aid any possible research in the future with this program, each subroutine will be listed with a brief explanation of the calling arguments.

1. Subroutine VERTMO

U = Zonal component of the horizontal wind in m sec^{-1} .

V = Meridional component of the horizontal wind in m sec^{-1} .

OMEGA = ω = Vertical component of the wind in (x, y, p, t) coordinate system in mb sec^{-1} .

NI = Number of grid points along x-axis.

NJ = Number of grid points along y-axis.

NK = Number of levels in the vertical.

DIST = Distance between grid points in meters.

DSI = The negative of the divergence field for 1000 mb = the sum of the divergence from 950 mb to 150 mb.

2. Subroutine FORFUN

VORT = Vertical component of relative vorticity.

F = Field for forcing function.

FMAP = Map factor (mean).

FBAR = Coriolis parameter (mean).

DIST = Distance between grid points in meters.

DT = Forward time step in seconds.

USPD = Zonal component of the horizontal wind in m sec^{-1} .

VSPD = Meridional component of the horizontal wind in m sec^{-1} .

K = The level the main part of the program is at when this subroutine is called.

IZ = Counter to determine at what time step (1-30) the main program is.

PSI = The ψ field in $\text{m}^2 \text{sec}^{-1}$.

OMEGA = ω = Vertical component of the wind in (x, y, p, t) coordinate system in mb sec^{-1} .

3. Subroutine RELAXI

M = Number from 2, NI = a number one greater than what is needed in DO LOOP.

N = Same as M except 2, NJ.

F = $\Delta\psi$ field for relaxation of balance equation.

A = ψ field in $\text{m}^2 \text{sec}^{-1}$.

B = Z field in meters.

L = Counter to keep track of number of points which have converged.

EPS = Epsilon = convergence criteria.

ALPHA = Over-relaxation constant.

FMAP = Map factor (mean).

FBAR = Coriolis parameter (mean).

DIST = Distance between grid points in meters.

IL = Lower point for DO LOOP in x direction.

JL = Lower point for DO LOOP in y direction.

LM = Iteration number from main program.

U = Zonal component of the horizontal wind in m sec^{-1} .

V = Meridional component of the horizontal wind in m sec^{-1} .

K = Level in the vertical from main program.

G = Gravity.

4. Subroutine RELAX

Same as in RELAXI except:

F = Forcing function.

R = Residual field.

ICOUNT = L.

I1 = IL.

J1 = JL.

5. Subroutine SMOOTH

D = Field to be smoothed.

RF = Smoothing factor (<1.0).

IPASS = Number of smoothing passes to make through the grid.

IL, JL = Lower point for DO LOOP.

IH, JH = Upper point for DO LOOP.

6. Subroutine PROG

TEMP = Temperature.

RH = Relative humidity.

DIR = Wind directions.

SPD = Wind speeds.

USPD = Zonal component of SPD in m sec^{-1} .

VSPD = Meridional component of SPD in m sec^{-1} .

DIST = Distance between grid points in meters.

NK = Number of levels in the vertical.

NI = Number of grid points on x-axis.

NJ = Number of grid points on y-axis.

NI1 = NI-1.

NJ1 = NJ-1.

OMEGA = ω .

Z = Height field.

7. Subroutine SETUP

TEMP = Temperature.

RH = Relative humidity.

Z = Height field.

IL = A number from 1 to NI, usually 1.

IH = A number from 1 to NI, usually NI.

JL = A number from 1 to NJ, usually 1.

JH = A number from 1 to NJ, usually NJ.

8. Subroutine CLOUDS.

This is the modified version of the W.D. model (Weinstein and Davis, 1967).

PDATA = Pressure levels.

ZDATA = Height field.

RHDATA = Relative humidity field.

ZZ1 = Height of cloud base.

DZ = Vertical increment (meters) to be used.

NCALL = Number of times this subroutine called. Used as
variable to bypass a READ statement.

IPT = I coordinate for set of data coming in.

JPT = J coordinate for set of data coming in.

9. Calling arguments of GRPHCL and GRAPH are all parameters which have been computed or set in subroutine CLOUDS.

10. Subroutine METMAP

This subroutine was not written by the author, but is part of the Naval Postgraduate School Computer Facility Library and has been included for information. It is simply a shading routine that will print out a field of data with a 0.5 inch grid spacing and contour the field.

Y = Two-dimensional field to be contoured.

N = Number of rows I in the array to be contoured.

M = Number of columns J in the array to be contoured.

T = Title for printout . Up to 96 columns of alpha-numeric information.

BND = Bandwidth desired for contouring.

AZ = Scaling constant; that is, each value will be multiplied by AZ before printing it out. Only the first three numbers to the right of the decimal point are on the output, so AZ can be used to control which numbers are to be printed out.

BZ = Additive constant. This is useful when working with D-values. Usually, however, it is 0.0. If used, this constant will be added to each value of Y before printing it out.

AMIN = Minimum value for subroutine to start contouring.

IJT = 0 means subroutine will compute the minimum value and start contouring at that value. A value must still be specified for AMIN but it will not be used.

ICON = 1 if contouring desired.

= 0 if no contouring desired.

When working with this subroutine, it must be remembered that the grid point $I = 1, J = 1$, is the upper left corner of the grid, and I increases downward while J increases to the right. If an array is defined with the lower left corner as the point (1, 1), the following set of FORTRAN statements will get the array in the proper order for the subroutine:

```
DO 5 I=1, 24
```

```
DO 5 J=1, 24
```

```
5 DUMMY (J,I)=Z(I,24-J+1)
```

The field called DUMMY is then taken into METMAP as Y.


```

C      FI=3.141592
C      DIST=1.853E3*5.0
C      DIST=DY=DX=GRID DISTANCE IN METERS
C      NK = 7
C      NK = NUMBER OF LEVELS IN THE VERTICAL
C      NK1=NK-1
C      NK2=NK-2
C      NI=24
C      NI = NUMBER OF GRID POINTS ALONG THE I-AXIS
C      NI1=NI-1
C      NI2=NI-2
C      NJ=24
C      NJ = NUMBER OF GRID POINTS ALONG THE J-AXIS
C      NJ1=NJ-1
C      NJ2=NJ-2
C      PHI=35.0*PI/180.
C      TWOPHI=2.0*PHI
C      FMAP=1.86603/(1.0+SIN(PHI))
C      G=9.780356*(1.0+.0052885*(SIN(PHI))**2-.0000059*(SIN(TWOPHI))**2)
C      G=GRAVITY FMAP=MAP FACTOR FOR 35 DEGREES LATITUDE
C      PEAD(5,6) ((( Z(I,J,K),I=1,NI),J=1,NJ),K=1,NK2)
C      PEAD(5,5) ((( Z(I,J,K),I=1,NI),J=1,NJ),K=NK1,NK)
C      5 FORMAT(12F6.0)
C      6 PEAD(5,6) ((( TEMP(I,J,K),I=1,NI),J=1,NJ),K=1,NK)
C      PEAD(5,6) ((( RH(I,J,K),I=1,NI),J=1,NJ),K=1,NK)
C      PEAD(5,6) ((( DIR(I,J,K),I=1,NI),J=1,NJ),K=1,NK)
C      PEAD(5,6) ((( SPD(I,J,K),I=1,NI),J=1,NJ),K=1,NK)
C      DO 8 K=1,NK
C      DO 8 I=1,NI
C      DO 8 J=1,NJ
C      DSI(I,J)=0.0
C      USPD(I,J,K) = -SPD(I,J,K)*SIN(DIR(I,J,K)/57.2958)
C      VSPD(I,J,K) = -SPD(I,J,K)*COS(DIR(I,J,K)/57.2958)
C      8 CONTINUE
C      NOW HAVE ALL FIELDS READ IN FOR NK LEVELS IN THE VERTICAL
C      DO 30 K=1,NK
C      DO 20 I=2,NI1
C      DO 20 J=2,NJ1
C      DUMMY(I,J)=1.0/(2.0*DIST))* (USPD(I+1,J,K)-USPD(I-1,J,K))+
C      IVSPD(I,J+1,K)-VSPD(I,J-1,K))
C      HEPE{ DUMMY USED TO REPRESENT DIVERGENCE FIELDS
C      DSI(I,J)=DSI(I,J)+DUMMY(I,J)
C      HERE, DSI USED TO REPRESENT THE SUM OF THE DIVERGENCE FIELDS IN A
C      COLUMN FROM 950MB TO 150MB
C      20 CONTINUE
C      30 CONTINUE
C      NUMB=NI2*NJ2

```

```

MN00490
MN00500
MN00510
MN00520
MN00530
MN00540
MN00550
MN00560
MN00570
MN00580
MN00590
MN00600
MN00610
MN00620
MN00630
MN00640
MN00650
MN00660
MN00670
MN00680
MN00690
MN00700
MN00710
MN00720
MN00730
MN00740
MN00750
MN00760
MN00770
MN00780
MN00790
MN00800
MN00810
MN00820
MN00830
MN00840
MN00850
MN00860
MN00870
MN00880
MN00890
MN00900
MN00910
MN00920
MN00930
MN00940
MN00950
MN00960

```

```

C NUMB = NUMBER OF POINTS TO BE RELAXED TO GET PSI FIELD
C EPS=58612.
C EPS IS THE CONVERGENCE CRITERIA FOR RELAXATION
C TIME=17.0
C ITM=1
C ITM WILL BE USED AS SUBSCRIPT FOR TEVEL
C TITLE IS TITLE
C TITLE IS TITLE INFORMATION FOR METMAP OF HEIGHTS
35 FORMAT(20A4, T5
C T5 IS TITLE INFORMATION FOR METMAP OF OMEGAS
C CALL VERTMO(USPD,VSPD,OMEGA,NI,NJ,NK,DIST,DSI)
CBND=40./3600.0
C CBND = BANDWIDTH FOR CONTOURING OF OMEGA FIELD IN METMAP
41 DO 300 K=1,NK
C THIS DO LOOP OBTAINS 1.5 HOUR FORECASTS FOR PSI FOR ALL SEVEN LEVELS
C TITLE(7) = BEVEL(K)
T5(7)=BEVEL(K)
C TITLE(23)=TEVEL(ITM)
C NEXT, CONVERT OMEGA TO 2-D FIELD FOR METMAP
DO 42 I=1,22
DO 42 J=1,22
DI(J,I)=OMEGA(I+1,24-J,K)
42 CONTINUE
C NOW GET PSI FIELD FIRST GUESS FROM THE INITIAL OBSERVED HEIGHTS
DO 55 I=1,NI
DO 55 J=1,NJ
DUMMY(J,I)=Z(I,NJ-J+1,K)
C DUMMY IS USED TO HAVE A 2-D FIELD TO TAKE INTO METMAP
SIOLD(I,J)=(G/FBAR)*Z(I,J,K)
SINew(I,J)=SIOLD(I,J)
DSI(I,J)=Z(I,J,K)
C DSI IS USED TO HAVE A 2-D FIELD TO TAKE INTO RELAXI
SIPROG(I,J)=SIOLD(I,J)
55 CONTINUE
C NOW USE METMAP TO CONTOUR THE ORIGINAL HEIGHT FIELD
C CALL METMAP(DUMMY,24,24,TITLE,BND(K),.001,0.0,ZBAR(K),0,1)
C NOW RELAX THE FIRST GUESS PSI FIELD
DO 60 ITER=1,200
L=0
CALL RELAXI(24,24,F,SIOLD,DSI,L,EPS,1.4,FMAP,FBAR,DIST,2,2,ITER,
1 USPD,VSPD,K,G)
IF(L.GE.NUMB) GO TO 64
60 CONTINUE
62 WRITE(6,62)
62 FORMAT(2X,'200 PASSES AND NO CONVERGENCE')
GO TO 999

```

MNO0970
MNO0980
MNO0990
MNO1000
MNO1010
MNO1020
MNO1030
MNO1035
MNO1040
MNO1050
MNO1055
MNO1060
MNO1070
MNO1080
MNO1090
MNO1100
MNO1110
MNO1120
MNO1130
MNO1135
MNO1140
MNO1150
MNO1160
MNO1170
MNO1180
MNO1190
MNO1200
MNO1210
MNO1220
MNO1230
MNO1240
MNO1250
MNO1260
MNO1270
MNO1280
MNO1290
MNO1295
MNO1300
MNO1310
MNO1320
MNO1330
MNO1340
MNO1350
MNO1360
MNO1370
MNO1380
MNO1390
MNO1400


```

64 WRITE(6,65) REVEL(K),L,ITER,TIME
65 FORMAT(11,' AT ',A4,' MB ',I5,' POINTS WERE WITHIN LIMITS (SHO
1ULD BE 484) ON PASS NUMBER',I5,' AT TIME=',F7.2,/)
C NOW SET UP THE EQUATIONS WHICH LEAD TO THE PROGGED PSI FIELD
NUMB=(NI-4)*(NJ-4)
IZ=1
C IZ = COUNTER WHICH REPRESENTS NUMBER OF TIME STEPS
DT=180.
C DT = FORWARD TIME STEP = .050 HOURS =180 SECONDS
C NOW MAKE FIRST GUESS AT THE DELTA PSI FIELD
DO 75 I=1,NI
DO 75 J=1,NJ
DSI(I,J)=0.0
75 CONTINUE
C NOW COMPUTE THE FORCING FUNCTION FOR THE PROG
CALL FORFUN(VORT,F,FMAP,FBAR,DIST,DT,USPD,VSPD,K,IZ,SIOLD,OMEGA)
C NOW RELAX THE DELTA PSI FIELD FOR FIRST TIME STEP
DO 90 ITER=1,100
L=0
CALL RELAX(NI,NJ,RESID,F,DSI,L,EPS,ALPHA,FMAP,3,3)
IF(L,GE,NUMB) GO TO 100
90 CONTINUE
100 DO 110 I=3,NI2
DO 110 J=3,NJ2
SIPROG(I,J)=SIOLD(I,J)+DSI(I,J)
SINew(I,J)=SIPROG(I,J)
110 CONTINUE
TIME=TIME+DT/3600.
DO 210 IZ=2,30
C STARTED AT IZ=2 BECAUSE PROG FOR TIME=1 ALREADY OBTAINED ABOVE
30 TIME STEPS ARE NEEDED FOR A 1.5 HOUR PROG
C UTILIZING THE LATEST PROG,A NEW FORCING FUNCTION FIELD
IS NEEDED BEFORE OBTAINING A NEW DELTA PSI FIELD
CALL FORFUN(VORT,F,FMAP,FBAR,DIST,DT,USPD,VSPD,K,IZ,SINew,OMEGA)
DO 150 ITER=1,50
L=0
CALL RELAX(NI,NJ,RESID,F,DSI,L,EPS,ALPHA,FMAP,3,3)
IF(L,GE,NUMB) GO TO 180
150 CONTINUE
C NOW HAVE A NEW DELTA PSI FIELD FOR THE 'IZ,TH' TIME STEP
180 TIME=TIME+DT/3600.
DO 182 I=3,NI2
DO 182 J=3,NJ2
SIPROG(I,J)=SIOLD(I,J)+2.0*DSI(I,J)

```

```

MNC1410
MNC1420
MNC1430
MNC1440
MNC1450
MNC1460
MNC1470
MNC1480
MNC1490
MNC1500
MNC1510
MNC1520
MNC1530
MNC1540
MNC1550
MNC1560
MNC1570
MNC1580
MNC1590
MNC1600
MNC1610
MNC1620
MNC1630
MNC1640
MNC1650
MNC1660
MNC1670
MNC1680
MNC1690
MNC1700
MNC1710
MNC1720
MNC1730
MNC1740
MNC1750
MNC1760
MNC1770
MNC1780
MNC1790
MNC1800
MNC1810
MNC1820
MNC1830
MNC1840
MNC1850
MNC1860
MNC1870
MNC1880

```

```

182 CONTINUE
C IF (I2.EQ.10.OR.I2.EQ.20.OR.I2.EQ.30) GO TO 186
C THIS STEP SMOOTHS EVERY HALF HOUR (10 TIME STEPS) TO KEEP FROM
C HAVING TWO INDEPENDENT SOLUTIONS FROM LEAP-FROG SCHEME
GO TO 195
186 DO 188 I=3,N12
DO 188 J=3,NJ2
RESID(I,J)=SINew(I,J)+DSI(I,J)
C HERE RESID USED AS DUMMY VARIABLE FOR SIPROG
SIPROG(I,J)=(SIPROG(I,J)+RESID(I,J))/2.0
188 CONTINUE
190 CALL SMOOTH(SIPROG,.04,1,3,22,3,22)
195 DO 200 I=3,N12
DO 200 J=3,NJ2
SIOLD(I,J)=SINew(I,J)
SINew(I,J)=SIPROG(I,J)
200 CONTINUE
210 CONTINUE
ITM=ITM+1

C THIS COMPLETES THE PSI PROG FOR 1.5 HOURS FOR THE K*TH LEVEL
C MUST RECOVER Z FIELD EACH 1.5 HOURS IN ORDER TO USE CLOUD SUB.
C MAKE INITIAL Z-FIELD GUESS
DO 220 I=3,N12
DO 220 J=3,NJ2
Z(I,J,K)={FBAR/G)*SINew(I,J)
220 CONTINUE
C NEXT, SET THE BOUNDARY CONDITIONS
DO 225 J=1,NJ
Z(2,J,K)=Z(3,J,K)-Z(3,J,K))
Z(1,J,K)=Z(2,J,K)-Z(3,J,K)-Z(2,J,K))
Z(23,J,K)=Z(22,J,K)-Z(22,J,K))
Z(23,NJ,K)=Z(N11,J,K)-Z(N11-1,J,K)-Z(N11,J,K))
225 CONTINUE
DO 230 I=1,N1
Z(I,2,K)=Z(I,3,K)-Z(I,3,K))
Z(I,1,K)=Z(I,2,K)-Z(I,3,K)-Z(I,2,K))
Z(I,23,K)=Z(I,22,K)-Z(I,21,K)-Z(I,22,K))
Z(I,NJ,K)=Z(I,NJ1,K)-Z(I,NJ1-1,K)-Z(I,NJ1,K))
230 CONTINUE
C NOW RELAX THE PROGGED Z-FIELD
NUMB = N12*NJ2
EPS=0.5
C SET UP THE FORCING FUNCTION FIELD
DO 235 I=1,N1

```

```

MNO1890
MNO1900
MNO1910
MNO1920
MNO1930
MNO1940
MNO1950
MNO1960
MNO1970
MNO1980
MNO1990
MNO2000
MNO2010
MNO2020
MNO2030
MNO2040
MNO2050
MNO2060
MNO2065
MNO2070
MNO2080
MNO2090
MNO2100
MNO2110
MNO2120
MNO2130
MNO2140
MNO2150
MNO2160
MNO2170
MNO2180
MNO2190
MNO2200
MNO2210
MNO2220
MNO2230
MNO2240
MNO2250
MNO2260
MNO2270
MNO2280
MNO2290
MNO2300
MNO2310
MNO2320
MNO2330
MNO2340
MNO2350

```



```

C      DO 235 J=1,NJ
C      DSI(I,J) = Z(I,J,K)
C      DSI IS BEING USED FOR Z TO HAVE 2-D FIELD FOR RELAX
235  CONTINUE
      DO 240 I=2,N11
      DO 240 J=2,NJ1
      T1=SINEW(I+1,J)+SINEW(I-1,J)+SINEW(I,J+1)+SINEW(I,J-1)-4.0*
1  SINEW(I,J)
      T2=2.0/(4.0*FBAR)
      T3=(USPD(I+1,J,K)-USPD(I-1,J,K))*(VSPD(I,J+1,K)-VSPD(I,J-1,K))
      T4=(USPD(I+1,J,K)-USPD(I-1,J,K))*(USPD(I,J+1,K)-USPD(I,J-1,K))
      F(I,J)=(FBAR/G)*(T1+T2*(T3-T4))
240  CONTINUE
      DO 250 ITER=1,200
      L=0
      CALL RELAX(NI,NJ,RESID,F,DSI,L,EPS,ALPHA,FMAP,2,2)
      IF(L.GE.NUMB) GO TO 251
250  CONTINUE
251  DO 252 I=1,N1
      DO 252 J=1,NJ
      Z(I,J,K) = DSI(I,J)
252  CONTINUE
C      METMAP IS NEXT TO CONTOUR THE Z-FIELD
C
C      DO 268 I=1,N1
C      DO 268 J=1,NJ
C      DUMMY(J,I)=Z(I,NJ-J+1,K)
268  CONTINUE
      TITLE(23)=TEVEL(ITM)
      CALL METMAP(DUMMY,24,24,TITLE,BND(K),.001,0.0,ZBAR(K),0,1)
C      THIS COMPLETES THE PSI PROG FOR ONE LEVEL(K) AHEAD FOR 1.5 HOURS
C
C      TIME=TIME-1.5
C      ITM=ITM-1
C      EPS=58612.
300  CONTINUE
301  TIME=TIME+1.5
      ITM=ITM+1
C      THIS COMPLETES ALL THE LEVELS FOR A 1.5 HOUR PROG OF PSI
C      NOW PROG AHEAD TEMPERATURE AND RELATIVE HUMIDITY
C      CALL PROG(TEMP,RH,DIR,SPD,USPD,VSPD,DIST,NK,NI,NJ,N11,NJ1,OMEGA)
C      NOW SETUP THE PROGGED SOUNDING AND TAKE IT INTO SUBROUTINE CLOUDS
C      CALL SETUP(TEMP,RH,Z,1,NI,1,NJ)
C      IF(TIME.GT.18.4) GO TO 41
      GO TO 41
999  STOP

```

```

MN02360
MN02370
MN02380
MN02390
MN02395
MN02400
MN02405
MN02410
MN02415
MN02420
MN02425
MN02430
MN02440
MN02450
MN02460
MN02470
MN02480
MN02490
MN02500
MN02510
MN02520
MN02530
MN02540
MN02550
MN02560
MN02570
MN02580
MN02590
MN02600
MN02605
MN02610
MN02620
MN02630
MN02640
MN02650
MN02660
MN02670
MN02680
MN02690
MN02700
MN02710
MN02720
MN02730
MN02740
MN02750
MN02760
MN02770
MN02780

```

END

MN02 790

```

SUBROUTINE VERTMO(U,V,OMEGA,NI,NJ,NK,DIST,DSI)
DIMENSION U(24,24,8),V(24,24,8),OMEGA(24,24,8),DSI(24,24)

THIS IS A SUBROUTINE TO COMPUTE VERTICAL MOTION ACCORDING TO THE
CONTINUITY EQUATION
BOUNDARY CONDITIONS ARE THAT OMEGA AT 150MB = 0 AND OMEGA AT
1000MB = 0
DSI IS THE NEGATIVE OF THE DIVERGENCE FIELD FOR 1000MB.
THE REASONING FOR THIS IS THAT THE DIVERGENCE AT 1000MB IS SET EQUAL
TO THE NEGATIVE OF THE SUM OF THE DIVERGENCES IN THE COLUMN FROM
150MB DOWN THRU 950MB.
HERE, OMEGA VALUES ARE COMPUTED FROM 150 MB DOWN THRU 300 MB
AND ARE COMPUTED UPWARDS FROM 1000 MB THRU 700 MB.
THEN, OMEGA AT 500MB IS COMPUTED AS THE AVERAGE OF THE OMEGA VALUES
AT 300MB AND 700MB

      NI1 = NI - 1
      NJ1 = NJ - 1
      DO 5 I=1,NI
      DO 5 J=1,NJ
      OMEGA(I,J,1) = 0.0
      OMEGA(I,J,8) = 0.0
      5 CONTINUE
      DO 200 K=1,6
      FOR THIS DO LOOP, THE K VALUES DO NOT CORRESPOND TO SPECIFIC LEVELS
      IN THE VERTICAL
      DO 190 I=2,NI1
      DO 190 J=2,NJ1
      IF(K.EQ.1) GO TO 60
      IF(K.EQ.2) GO TO 50
      IF(K.EQ.3) GO TO 40
      IF(K.EQ.4) GO TO 10
      IF(K.EQ.5) GO TO 20
      IF(K.EQ.6) GO TO 30
      *****
      THE NEXT SET OF STATEMENTS COMPUTES OMEGA AT 200MB AS THE AVERAGE
      BETWEEN 175MB AND 225MB
      10 W175=OMEGA(I,J,8)-(25./(4.*DIST))*(U(I+1,J,7)-U(I-1,J,7))+
      1V(I,J+1,7)-V(I,J-1,7)+U(I-1,J,6)+V(I,J+1,6)-V(I,J-1,6))
      NOW HAVE OMEGA AT 175MB
      NEXT GET OMEGA AT 225MB, THEN AVERAGE IT WITH OMEGA AT 175MB

```

V000010
 V000020
 V000030
 V000040
 V000050
 V000060
 V000070
 V000080
 V000090
 V000100
 V000110
 V000120
 V000130
 V000140
 V000150
 V000160
 V000170
 V000180
 V000190
 V000200
 V000210
 V000220
 V000230
 V000240
 V000250
 V000260
 V000270
 V000280
 V000290
 V000300
 V000310
 V000320
 V000330
 V000340
 V000350
 V000360
 V000370
 V000380
 V000390
 V000400
 V000410
 V000420
 V000430

```

C      TO GET OMEGA AT 200MB
      W225=OMEGA(I,J,8)-(75./(4.*DIST))*(U(I+1,J,7)-U(I-1,J,7)+V(I,J+1,
17)-V(I,J-1,7))+U(I+1,J,5)-U(I-1,J,5)+V(I,J+1,5)-V(I,J-1,5))
      OMEGA(I,J,7) = (W175 + W225)/2.0
      OMEGA(I,J,7) IS NOW THE OMEGA FOR 200MB
      GO TO 190
C
C
C      *****
C      THE NEXT SET OF STATEMENTS COMPUTES OMEGA AT 300MB AS THE AVERAGE
C      BETWEEN 350MB AND 250MB
C      FIRST, CCMPUTE OMEGA AT 350MB
20  W350=OMEGA(I,J,7)-(150./(4.*DIST))*(U(I+1,J,6)-U(I-1,J,6)+V(I,J+1,
1J+1,6)-V(I,J-1,6))+U(I+1,J,4)-U(I-1,J,4)+V(I,J+1,4)-V(I,J-1,4))
C      NOW COMPUTE OMEGA AT 250MB, THEN AVERAGE IT WITH OMEGA AT
350MB TO GET OMEGA AT 300MB
      W250=OMEGA(I,J,7)-(50./(4.*DIST))*(U(I+1,J,6)-U(I-1,J,6)+V(I,J+1,
16)-V(I,J-1,6))+U(I+1,J,5)-U(I-1,J,5)+V(I,J+1,5)-V(I,J-1,5))
      OMEGA(I,J,6) = (W350 + W250)/2.0
      OMEGA(I,J,6) IS NOW THE OMEGA FOR 300MB
      GO TO 190
C
C
C      *****
C      THE NEXT SET OF STATEMENTS COMPUTES OMEGA AT 500MB AS THE VALUE
C      OBTAINED WHEN THE OMEGA VALUES OF 300MB AND 700MB ARE AVERAGED
30  OMEGA(I,J,5)=(OMEGA(I,J,6)+OMEGA(I,J,4))/2.0
      OMEGA(I,J,5) IS NOW THE OMEGA FOR 500MB
      GO TO 190
C
C
C      *****
C      THE NEXT SET OF STATEMENTS COMPUTES OMEGA AT 700MB AS THE
C      AVERAGE BETWEEN OMEGA AT 675MB AND OMEGA AT 725MB
40  W675=OMEGA(I,J,3)+(175./(4.*DIST))*(U(I+1,J,4)-U(I-1,J,4)+V(I,J+1,
1J+1,4)-V(I,J-1,4))+U(I+1,J,2)-U(I-1,J,2)+V(I,J+1,2)-V(I,J-1,2))
C      NOW COMPUTE OMEGA AT 725MB
      W725=OMEGA(I,J,2)+(225./(4.*DIST))*(U(I+1,J,4)-U(I-1,J,4)+V(I,J+1,
14)-V(I,J-1,4))+U(I+1,J,1)-U(I-1,J,1)+V(I,J+1,1)-V(I,J-1,1))
      OMEGA(I,J,4) = (W675 + W725)/2.0
      OMEGA(I,J,4) IS NOW OMEGA AT 700MB
      GO TO 190
C
C
C      *****
C      THE NEXT SET OF STATEMENTS CCMPUTES OMEGA AT 850MB AS THE
C      VALUE OBTAINED WHEN VERTICAL MOTIONS AT 825MB AND 900MB ARE AVERAGED
      USING A WEIGHTED TYPE MEAN

```

```

V000440
V000450
V000460
V000470
V000480
V000490
V000500
V000510
V000520
V000530
V000540
V000550
V000560
V000570
V000580
V000590
V000600
V000610
V000620
V000630
V000640
V000650
V000660
V000670
V000680
V000690
V000700
V000710
V000720
V000730
V000740
V000750
V000760
V000770
V000780
V000790
V000800
V000810
V000820
V000830
V000840
V000850
V000860
V000870
V000880
V000890
V000900
V000910

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```

00009200
00009300
00009400
00009500
00009600
00009700
00009800
00009900
00010000
00010100
00010200
00010300
00010400
00010500
00010600
00010700
00010800
00010900
00011000
00011100
00011200
00011300
00011400
00011500
00011600
00011700
00011800
00011900
00012000
00012100
00012200
00012300
00012400
00012500
00012600
00012700
00012800
00012900
00013000
00013100
00013200
00013300
00013400
00013500
00013600

50 W825=OMEGA(I,J,2)+(125./(4.*DIST))*((U(I+1,J,3)-U(I-1,J,3))+
1V(I,J+1,3))-V(I,J-1,3)+U(I+1,J,1))-U(I-1,J,1))+V(I,J+1,1))-V(I,J-1,1))
W900=OMEGA(I,J,2)+{ 50./(4.*DIST))*((U(I+1,J,2)-U(I-1,J,2))+
1V(I,J+1,2))-V(I,J-1,2)+U(I+1,J,1))-U(I-1,J,1))+V(I,J+1,1))-V(I,J-1,1))
OMEGA(I,J,3)}3.0
OMEGA(I,J,3) IS NOW OMEGA AT 850MB
GO TO 190

*****
THE NEXT SET OF STATEMENTS COMPUTES OMEGA AT 950MB AS THE
AVERAGE OF THE OMEGA VALUES AT 925MB AND 975MB
60 W925=OMEGA(I,J,1)+(75.0/2.0)*((U(I+1,J,2)-U(I-1,J,2))+
1V(I,J+1,2))-V(I,J-1,2))/(2.0*DIST))-DSI(I,J))
W975= 0.0+12.5*((U(I+1,J,1))-U(I-1,J,1))+V(I,J+1,1))-V(I,J-1,1))/
1(2.0*DIST))-DSI(I,J))
NOW AVERAGE W925 AND W975 TO GET OMEGA AT 950MB
OMEGA(I,J,2) = (W925 + W975)/2.0
OMEGA(I,J,2) IS NOW OMEGA AT 950MB

*****
190 CONTINUE
200 CONTINUE

THE THIRD SUBSCRIPT ON OMEGA IS OFF ONE FROM THE THIRD
SUBSCRIPT OF THE OTHER PARAMETERS, SO SINCE ALL OMEGAS HAVE NOW
BEEN COMPUTED, IT IS NECESSARY TO REVERT THE SUBSCRIPTS BACK SO
THEY WILL MATCH

DO 250 K=1,7
DO 225 I=2,N11
DO 225 J=2,NJ1
OMEGA(I,J,K) = OMEGA(I,J,K+1)*3600.0
225 CONTINUE
ALL OMEGAS ARE NOW IN UNITS OF MILLIBARS PER HOUR
DO 245 I=2,N11
DO 245 J=2,NJ1
OMEGA(I,J,K) = OMEGA(I,J,K)/3600.0
245 CONTINUE
OMEGA IS NOW IN MILLIBARS PER SECOND
250 CONTINUE
RETURN
END

```



```

C
C
C
C
SUBROUTINE FORFUN(VORT,F,FMAP,FBAR,DIST,DT,USPD,VSPD,K,IZ,PSI,
1 OMEGA)
DIMENSION VORT(24,24,8),F(24,24),USPD(24,24,8),VSPD(24,24,8),
1 OMEGA(24,24,8),PSI(24,24)
THIS SUBROUTINE COMPUTES VORTICITY FIELDS THE FIRST TIME IT IS
CALLED AND COMPUTES THE FORCING FUNCTION EVERY TIME.
IF(K.EQ.1) DP=100.0
IF(K.EQ.2) DP=250.0
IF(K.EQ.3) DP=350.0
IF(K.EQ.4) DP=400.0
IF(K.EQ.5) DP=300.0
IF(K.EQ.6) DP=150.0
IF(K.EQ.7) DP=50.0
IF(IZ.GE.2) GO TO 215
DO 210 KK=1,7
DO 200 I=2,23
DO 200 J=2,23
VORT(I,J,KK)=(FMAP/(2.0*DIST))*(VSPD(I+1,J,KK)-VSPD(I-1,J,KK))-
1 USPD(I,J,KK)+USPD(I,J,KK)
200 CONTINUE
210 CONTINUE
215 KK=K
C STATEMENT 215 IS MERELY A DUMMY STATEMENT
DO 300 I=3,22
DO 300 J=3,22
SIVORT=((FMAP**2)/(DIST**2))*(PSI(I+1,J)+PSI(I-1,J)+PSI(I,J+1)+
1 PSI(I,J-1))-4.0*PSI(I,J))
SIVORT IS THE LAPLACIAN OF THE STREAM FIELD AND IS USED
FOR THE VORTICITY FIELD IN THE DIVERGENCE TERM OF THE VORTICITY EON.
IF(K.EQ.1) GO TO 250
IF(K.EQ.7) GO TO 275
A=OMEGA(I,J,K)*(VORT(I,J,K+1)-VORT(I,J,K-1))/DP
R=-((FMAP**2.0)/(4.0*DIST**2.0))*((PSI(I,J+1)-PSI(I,J-1))*
1 (VORT(I+1,J,K)-VORT(I-1,J,K))-(PSI(I+1,J)-PSI(I-1,J))*
2 (VORT(I,J+1,K)-VORT(I,J-1,K)))
C=(SIVORT+FBAR)*((USPD(I+1,J,K)-USPD(I-1,J,K)+VSPD(I,J+1,K)-
1 VSPD(I,J-1,K))/(2.0*DIST))
D=((OMEGA(I+1,J,K)-OMEGA(I-1,J,K))/(2.0*DIST))*((VSPD(I,J,K-1)-
1 VSPD(I,J,K+1))/DP)
E=-((OMEGA(I,J+1,K)-OMEGA(I,J-1,K))/(2.0*DIST))*((USPD(I,J,K-1)-
1 USPD(I,J,K+1))/DP)
F(I,J)=((-1.0*DIST**2.0*DT)/(FMAP**2.0))*(A+B+C+D+E)
GO TO 300
250 A=OMEGA(I,J,K)*(VORT(I,J,K)-VORT(I,J,K+1))/DP
R=-((FMAP**2.0)/(4.0*DIST**2.0))*((PSI(I,J+1)-PSI(I,J-1))*
1 (VORT(I+1,J,K)-VORT(I-1,J,K))-(PSI(I+1,J)-PSI(I-1,J))*

```

FN00010
FN00020
FN00030
FN00040
FN00050
FN00060
FN00065
FN00070
FN00080
FN00090
FN00100
FN00110
FN00120
FN00130
FN00140
FN00150
FN00160
FN00170
FN00180
FN00190
FN00200
FN00210
FN00220
FN00230
FN00240
FN00250
FN00260
FN00270
FN00280
FN00290
FN00300
FN00310
FN00320
FN00330
FN00340
FN00350
FN00360
FN00370
FN00380
FN00390
FN00400
FN00410
FN00420
FN00430
FN00440
FN00450
FN00460
FN00470

```

2 (VORT(I,J+1,K)-VORT(I,J-1,K))
C=(SI VORT(I,J+1,K)+FBAR)*{(USPD(I+1,J,K)-USPD(I-1,J,K))+VSPD(I,J+1,K)-
1 VSPD(I,J-1,K))/(2.0*DIST)}
D=((OMEGA(I+1,J,K)-OMEGA(I-1,J,K))/(2.0*DIST))*{(VSPD(I,J,K)-
1 VSPD(I,J,K+1))/DP}
E=-((OMEGA(I,J+1,K)-OMEGA(I,J-1,K))/(2.0*DIST))*{(USPD(I,J,K)-
1 USPD(I,J,K+1))/DP}
F(I,J)=((-1.0*DIST**2.0*DT)/(FMAP**2.0))*(A+B+C+D+E)
GO TO 300
275 A=OMEGA(I,J,K)*(VORT(I,J,K-1)-VORT(I,J,K))/DP
B=-((FMAP**2.0)/(4.0*DIST**2.0))*{(PSI(I,J+1)-PSI(I,J-1))*
1 (VORT(I+1,J,K)-VORT(I-1,J,K))-((PSI(I+1,J)-PSI(I-1,J))*
2 (VORT(I,J+1,K)-VORT(I,J-1,K)))}
C=(SI VORT(I,J+1,K)+FBAR)*{(USPD(I+1,J,K)-USPD(I-1,J,K))+VSPD(I,J+1,K)-
1 VSPD(I,J-1,K))/(2.0*DIST)}
D=((OMEGA(I+1,J,K)-OMEGA(I-1,J,K))/(2.0*DIST))*{(VSPD(I,J,K-1)-
1 VSPD(I,J,K))/DP}
E=-((OMEGA(I,J+1,K)-OMEGA(I,J-1,K))/(2.0*DIST))*{(USPD(I,J,K-1)-
1 USPD(I,J,K))/DP}
F(I,J)=((-1.0*DIST**2.0*DT)/(FMAP**2.0))*(A+B+C+D+E)
300 CONTINUE
RETURN
END

```

FN00480
FN00490
FN00500
FN00510
FN00520
FN00530
FN00540
FN00550
FN00560
FN00570
FN00580
FN00590
FN00600
FN00610
FN00620
FN00630
FN00640
FN00650
FN00660
FN00670
FN00680
FN00690
FN00700

```

SUBROUTINE RELAXI(M,N,F,A,B,L,EPS,ALPHA,FMAP,FBAR,DIST,IL,JL,LM,
1 U,V,K,G)
DIMENSION A(24,24),B(24,24),F(24,24),R(24,24),U(24,24,8),
1 V(24,24,8)

```

RI00010
RI00020
RI00030
RI00040
RI00050
RI00060
RI00070
RI00080
RI00090
RI00100
RI00110
RI00120
RI00130
RI00140
RI00150
RI00160
RI00170
RI00180
RI00190
RI00200
RI00210

```

L=NUMBER OF POINTS CONVERGED
A=PSI FIELD
B=HEIGHT(Z) FIELD
R=RESIDUAL
F=FORCING FUNCTION
IH=M-1
JH=N-1
DO 100 I=IL,IH
DO 100 J=JL,JH
W=(U(I+1,J,K)-U(I-1,J,K))*{(V(I,J+1,K)-V(I,J-1,K))
X=(V(I+1,J,K)-V(I-1,J,K))*{(U(I,J+1,K)-U(I,J-1,K))
Z=(G/FBAR)*{(B(I+1,J)+B(I-1,J))+B(I,J+1)+B(I,J-1))-4.0*B(I,J)}
Y=A(I+1,J)+A(I-1,J)+A(I,J+1)+A(I,J-1)-4.0*A(I,J)
R(I,J)=Y+(2.0/(4.0*FBAR))*{(W-X)-Z
IF(ABS(R(I,J)).LT.EPS) L=L+1
DELTSI=ALPHA*R(I,J)/4.0

```

CC
CC
CC
CC
CC

```

      A(I,J)=A(I,J)+DELTSI
      F(I,J)=DELTSI
100  CONTINUE
      RETURN
      END

```

```

RI00220
RI00230
RI00240
RI00250
RI00260

```

```

      SUBROUTINE RELAX(M,N,R,F,A,ICOUNT,EPS,ALPHA,FMAP,IL,J1)
      DIMENSION R(24,24),F(24,24),A(24,24)
      C  ICOUNT = NUMBER OF POINTS CONVERGED
      IHIGH=M-1
      JHIGH=N-1
      DO 100 I=1,IHIGH
      DO 100 J=1,JHIGH
      R(I,J)=A(I+1,J)+A(I-1,J)+A(I,J+1)+A(I,J-1)-4.0*A(I,J)-F(I,J)
      IF(ABS(R(I,J)).LT.EPS) ICOUNT=ICOUNT+1
      A(I,J)=A(I,J)+(ALPHA/4.0)*R(I,J)
100  CONTINUE
      RETURN
      END

```

```

RX00010
RX00020
RX00030
RX00040
RX00050
RX00060
RX00070
RX00080
RX00090
RX00100
RX00110
RX00120
RX00130

```

```

      SUBROUTINE SMOOTH(D,RF,IPASS,IL,IH,JL,JH)
      DIMENSION D(24,24)
      DO 450 IK=1,IPASS
      DO 400 I=IL,IH
      DO 400 J=JL,JH
      PL=D(I,J+1)+D(I,J-1)+D(I+1,J)+D(I-1,J)-4.0*D(I,J)
      D(I,J)=D(I,J)+RF*PL
400  CONTINUE
450  CONTINUE
      RETURN
      END

```

```

SH00010
SH00020
SH00030
SH00040
SH00050
SH00060
SH00070
SH00080
SH00090
SH00100
SH00110

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SUBROUTINE PROG(TEMP,RH,DIR,SPD,USPD,VSPD,DIST,NK,NI,NJ,NJ1,NJL,
1 OMEGA)
C THIS SUBROUTINE OBTAINS FORECASTS OF TEMP. AND RELATIVE HUMIDITY
DIMENSION TEMP(24,24,8),RH(24,24,8),DIR(24,24,8),SPD(24,24,8),
1 USPD(24,24,8),VSPD(24,24,8),DTEMP(24,24),DRH(24,24),
2 OMEGA(24,24,8),DP(7)
DATA DP/-5.0E1,-1.0E2,-1.5E2,-2.0E2,-2.0E2,-1.0E2,-5.0E1/
C THIS SUBROUTINE IS DESIGNED TO PROG AHEAD BY TWO METHODS.
THE FIRST POSSIBILITY IS TO SET THE LOCAL TERM EQUAL TO THE
NEGATIVE OF THE HORIZONTAL ADVECTION TERM AND THE NEGATIVE OF
THE VERTICAL ADVECTION TERM. THE SECOND POSSIBILITY IS TO ONLY
USE THE NEGATIVE OF THE HORIZONTAL ADVECTION TERM. IN EITHER
CASE, ONLY 50% OF THE ADVECTION IS USED. TO INCLUDE THE VERTICAL
ADVECTION TERM, TWO CARDS ARE NEEDED BETWEEN THE SECOND DO 50 AND
THE STATEMENT JUST AFTER IT.
C
C DELTAT = 90 * 60.
C DELTAT = TIME STEP = 1.5 HOURS
DO 75 K=1,NK
DO 50 I=2,NJ1
DO 50 J=2,NJL
DTEMP(I,J)=-DELTAT*(USPD(I,J,K)*((TEMP(I+1,J,K)-TEMP(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)*((TEMP(I,J+1,K)-TEMP(I,J-1,K))/(2.*DIST)))
DRH(I,J)=-DELTAT*(USPD(I,J,K)*((RH(I+1,J,K)-RH(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)*((RH(I,J+1,K)-RH(I,J-1,K))/(2.*DIST)))
GO TO 50
20 DTEMP(I,J)=-DELTAT*(USPD(I,J,K)*((TEMP(I+1,J,K)-TEMP(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)*((TEMP(I,J+1,K)-TEMP(I,J-1,K))/(2.*DIST)))
2 OMEGA(I,J,K)*((TEMP(I,J,K)-TEMP(I,J,K))/DP(K+1))
DRH(I,J)=-DELTAT*(USPD(I,J,K)*((RH(I+1,J,K)-RH(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)*((RH(I,J+1,K)-RH(I,J-1,K))/(2.*DIST)))
2 OMEGA(I,J,K)*((RH(I,J,K)-RH(I,J,K))/DP(K+1))
GO TO 50
30 DTEMP(I,J)=-DELTAT*(USPD(I,J,K)*((TEMP(I+1,J,K)-TEMP(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)/2.0)*((TEMP(I,J,K)-TEMP(I,J,K))/(2.*DIST)))
3 OMEGA(I,J,K)-TEMP(I,J,K))/DP(K+1))
DRH(I,J)=-DELTAT*(USPD(I,J,K)*((RH(I+1,J,K)-RH(I-1,J,K))/
1 (2.*DIST)))+VSPD(I,J,K)/2.0)*((RH(I,J,K)-RH(I,J,K))/(2.*DIST)))
2 OMEGA(I,J,K)-RH(I,J,K))/DP(K+1))
3 OMEGA(I,J,K)-RH(I,J,K))/DP(K+1))
50 CONTINUE
IF(K.EQ.1) GO TO 69
DO 55 I=2,NJ1
DO 55 J=2,NJL
TEMP(I,J,K-1)=TEMP(I,J,8)
RH(I,J,K-1)=RH(I,J,8)

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```

55 CONTINUE
69 DO 70 I=2,NJ1
   DO 70 J=2,NJ1
     TEMP(I,J,8) = TEMP(I,J,K) +.5* DTEMP(I,J)
     RH(I,J,8) = RH(I,J,K) +.5*DRH(I,J)
     IF(RH(I,J,8).GT.100.) RH(I,J,8) = 100.0
     IF(RH(I,J,8).LT.10.) RH(I,J,8) = 10.0
70 CONTINUE
75 CONTINUE
   DO 76 I=2,NJ1
     DO 76 J=2,NJ1
       TEMP(I,J,7)=TEMP(I,J,8)
       RH(I,J,7)=RH(I,J,8)
76 CONTINUE
   DO 100 K=1,NK
     NEXT, SET THE BOUNDARY CONDITIONS
     DO 80 J=1,NJ
       TEMP(1,J,K)=TEMP(2,J,K)-(TEMP(3,J,K)-TEMP(2,J,K))
       TEMP(NI,J,K)=TEMP(NI-1,J,K)-(TEMP(NI,J,K)-TEMP(NI-1,J,K))
       RH(1,J,K)=RH(2,J,K)-(RH(3,J,K)-RH(2,J,K))
       IF(RH(1,J,K).GT.100.) RH(1,J,K)=100.0
       IF(RH(1,J,K).LT.10.) RH(1,J,K)=10.0
       RH(NI,J,K)=RH(NI-1,J,K)-(RH(NI,J,K)-RH(NI-1,J,K))
       IF(RH(NI,J,K).GT.100.) RH(NI,J,K)=100.0
       IF(RH(NI,J,K).LT.10.) RH(NI,J,K)=10.0
80 CONTINUE
     DO 85 I=1,NI
       TEMP(I,1,K)=TEMP(I,2,K)-(TEMP(I,3,K)-TEMP(I,2,K))
       TEMP(I,NJ,K)=TEMP(I,NJ-1,K)-(TEMP(I,NJ,K)-TEMP(I,NJ-1,K))
       RH(I,1,K)=RH(I,2,K)-(RH(I,3,K)-RH(I,2,K))
       IF(RH(I,1,K).GT.100.) RH(I,1,K)=100.0
       IF(RH(I,1,K).LT.10.) RH(I,1,K)=10.0
       RH(I,NJ,K)=RH(I,NJ-1,K)-(RH(I,NJ,K)-RH(I,NJ-1,K))
       IF(RH(I,NJ,K).GT.100.) RH(I,NJ,K)=100.0
       IF(RH(I,NJ,K).LT.10.) RH(I,NJ,K)=10.0
85 CONTINUE
100 CONTINUE
C THIS NOW COMPLETES THE PROG FOR BOTH PARAMETERS FOR 1.5 HOURS
C
C RETURN
C END

```

PG00490
 PG00500
 PG00510
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 PG00530
 PG00540
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 PG00580
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 PG00600
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 PG00690
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 PG00800
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 PG00880
 PG00890
 PG00900

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C      SUBROUTINE SETUP(TEMP,RH,Z,IL,IH,JL,JH)
C      THIS SUBROUTINE TAKES THE PROGRESSED Z-FIELD, CLOUDS, TEMP., AND REL. HUM.,
C      DETERMINES CLOUD BASE AND HAS SUBROUTINE CLOUDS PLOT THE RESULTS OF
C      ANY CONVECTIVE MOTION FOR EACH GRID POINT AT THE SPECIFIED INTERVALS
C      OF TIME AS SPECIFIED IN THE MAIN PROGRAM
      DIMENSION Z(24,24,8),TEMP(24,24,8),RH(24,24,8),P(8),A(8),B(8),
      1C(8)
      NCALL=1
      P(1)=950.
      P(2)=850.
      P(3)=700.
      P(4)=500.
      P(5)=300.
      P(6)=200.
      P(7)=150.
      P(8)=0.
C      COMPUTE CLOUD BASE, AND CALL CLOUDS FOR EACH GRID POINT
      DO 99 I=IL,IH
      DO 99 J=JL,JH
      10 IF(I.EQ.10.AND.J.EQ.1) GO TO 30
      11 IF(I.EQ.11.AND.J.EQ.3) GO TO 30
      12 IF(I.EQ.11.AND.J.EQ.9) GO TO 30
      13 IF(I.EQ.12.AND.J.EQ.9) GO TO 30
      14 IF(I.EQ.13.AND.J.EQ.14) GO TO 30
      15 IF(I.EQ.15.AND.J.EQ.17) GO TO 30
      16 IF(I.EQ.21.AND.J.EQ.18) GO TO 30
      17 IF(I.EQ.11.AND.J.EQ.23) GO TO 30
      18 GO TO 99
      19 DZ=200.0
      20 HEIGHT=Z(I,J,2)
      21 HEIGHT=THE HEIGHT OF CLOUD BASE IN METERS
      22 WRITE(6,35) TEMP(I,J,2),RH(I,J,2),Z(I,J,2)
      23 35 FORMAT(1,2X,'BEFORE LCL COMPUTED,TEMP(I,J,2)=',F10.1,3X,
      24 1,RH(I,J,2)=',F10.1,3X,'AND Z(I,J,2)=',F10.1,/)
C      NOW COMPUTE THE CLOUD STRUCTURE IN THE VERTICAL WITH W.D. MODEL
      DO 50 KK=1,8
      25 A(KK)=TEMP(I,J,KK)
      26 B(KK)=RH(I,J,KK)
      27 C(KK)=Z(I,J,KK)
      28 CONTINUE
      29 CALL CLOUDS(P,C,A,B,HEIGHT,DZ,NCALL,I,J)
      30 NCALL=NCALL+1
      31 CONTINUE
      32 99 RETURN
      33 END

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SP00010
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 SP00190
 SP00200
 SP00210
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 SP00240
 SP00250
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 SP00270
 SP00280
 SP00290
 SP00300
 SP00310
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 SP00470
 SP00480

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C SUBROUTINE CLOUDS(PDATA,ZDATA,TDATA,RHDATA,ZZ1 ,DZ,NCALL,IPT,JPT)
C
C SINGLE PRECISION
C
C W.D. PARAMETERIZED NUMERICAL MODEL CF CUMULUS CONVECTION
C MODIFIED AT NPGS
C
C DIMENSION P(2),T(2),X(2),Z(2),Q(2),QCL(2),QH(2),TVE(2),W(2),U(2),
C 1TH(2),RH(2),DEN(2),PE(200),TE(200),XE(200),ZE(200),UE(200),
C 2THE(200),RHE(200),AW(200),AQH(200),AQC(200),TC(200),AX(200),
C 3UPRAD(200),AO(3),RA(3),DUR(3),ITOP(3),NTF(3),SIZE(3),
C 4AREA(3),TEMP(3),TITLE(3),TITLE(19),PDATA(8),ZDATA(8),TDATA(8),RHDATA(8)
C AP=6958.9262
C BP=5.65567
C CX=59.01383
C
C AP,BP AND CX ARE CONSTANTS USED TO COMPUTE SATURATION VAPOR PRESSURE
C
C 100 READ(5,110) GO TO 111
C 110 FORMAT(19A4,I4)
C 111 CONST=273.16
C
C CONST IS CONVERSION FROM CELSIUS TO KELVIN DEGREES
C
C RD=287.04
C RV=461.5
C RATIO=RD/RV
C CP=1004.00
C
C RD,RV AND CP HAVE UNITS OF JOULES PER KILOGRAM PER DEGREE KELVIN
C
C THE R'S ARE GAS CONSTANTS FOR DRY AIR AND WATER VAPOR CP=SPEC. HEAT
C
C PI=3.14159265
C PHI=35.0
C PHIRAD= PHI*PI/180.
C TWOPHI=2.0*PHIRAD
C G=9.780356*(1.0+.0052885*( SIN(PHIRAD))**2-.0000059*( SIN(TWOPHI))
C 1*2)
C
C PHI IS MEAN LATITUDE FOR NSSL NETWORK G=GRAVITY
C
C CON1=-.62340
C CON2=621.7
C TNOT=233.16
C
C CON1=OPTIMUM DIFFERENCE OF SPECIFIC HEATS
C CON2= LATENT HEAT OF CONDENSATION IN CALORIES/GRAM AT 233 DEGREES
C
C 119 WRITE(6,120) TITLE,IPT,JPT
C 120 FORMAT( /,36X,19A4,/,37X,'GRID POINT IS I=',I4,4X,'J=',I4)
C
C ***** PART 1 INTERPOLATION *****
C

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CS00300
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CS00370
CS00380
CS00390
CS00400
CS00410
CS00420
CS00430
CS00440

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C
130 WRITE(6,130)
    FORMAT(//,60X,16HINITIAL SOUNDING,/)
    N=2
140 WRITE(6,140)
    FORMAT(20X,'PRESSURE',3X,'HEIGHT',5X,'TEMPERATURE',3X,'RELATIVE HU
1M,3X,'WIND SPEED',3X,'WIND DIRECTION',/,22X,'(MB)',4X,'(METERS)',.
26X,'(DEG C)',10X,'(%)',8X,'(M/SEC)',6X,'(DIRN. FROM)',/,)
    KK=2
    P(1)=PDATA(KK)
    T(1)=TDATA(KK)
    RH(1)=RHDATA(KK)
    U(1)=0.0
    TH(1)=0.0
    Z(1)=ZZ1
150 FORMAT (2F10.0,10X,4F10.0,F10.0)
    NDZ=DZ
    INDOC=200/NDZ
    TE(1)=T(1)+CONST
    PE(1)=P(1)
    IHT1=Z(1)
    RHE(1)=RH(1)/100.0
    ZE(1)=Z(1)
    UE(1)=U(1)
    THE(1)=TH(1)
    ILEVL=0
    KK = KK + 1
160 P(2)=PDATA(KK)
    T(2)=TDATA(KK)
    RH(2)=RHDATA(KK)
    U(2)=0.0
    TH(2)=0.0
170 FORMAT (2F10.0,10X,3F10.0)
    P(1),Z(1),T(1),RH(1),U(1),TH(1)
180 WRITE(6,180)
    FORMAT(19X,F8.2,F11.2,F14.2,F15.1,F13.2,F17.0)
190 IF(P(2)) 260,260,190
    Z(2)=Z(1)+ALOG(P(1)/P(2))*RD*(T(1)+T(2)+(2.0*CONST))/(2.0*G)
    XLNP=ALOG(P(2)/P(1))
    DTDZ=(T(2)-T(1))/XLNP
    DRHDP=(RH(2)-RH(1))/(XLNP*100.0)
    DUDP=(U(2)-U(1))/XLNP
    DTHDP=(TH(2)-TH(1))/XLNP
    DO 240 J=N,200
200 IF(ILEVL) 200,200,210
    PE(J)=PE(J-1)*EXP(-G*DZ/(RD*TE(J-1)))
    CP=ALOG(PE(J)/PE(J-1))
    TE(J)=TE(J-1)+(DTDZ*DP)
    RHE(J)=RHE(J-1)+(DRHDP*DP)

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00450
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 00880
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 00910
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 00920
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210 UE(J)=UE(J-1)+(DUDP*DP)
    THE(J)=THE(J-1)+(DTHDP*DP)
    GO TO 220
    PE(J)=P(1)* EXP(-G*(ZE(2)-Z(1)))/(RD*(T(1)+CONST)))
    DP=ALOG(PE(J)/P(1))
    TE(J)=T(1)+(DTPD*DP)+CONST
    RHE(J)=RH(1)/100.0+(DRHDP*DP)
    UE(J)=U(1)+(DUDP*DP)
    THE(J)=TH(1)+(DTHDP*DP)
    ILEV=0
    ZE(2)=ZE(1)+DZ
    IF(ZE(2)-Z(2)) 230,230,250
    ZE(1)=ZE(2)
    CONTINUE
    Z(1)=Z(2)
    P(1)=P(2)
    T(1)=T(2)
    U(1)=U(2)
    TH(1)=TH(2)
    RH(1)=RH(2)
    N=J
    ILEV=1
    KK=KK+1
    GO TO 160
260 IF(JN) 310,310,270
270 WRITE(6,280)
280 FORMAT(1,60X,'INTERPOLATED SOUNDING',//)
285 WRITE(6,140)
    FORMAT(20X,'PRESSURE',3X,'HEIGHT',3X,'TEMPERATURE',3X,'RELATIVE HU
1M',3X,'WIND SPEED',3X,'WIND DIRECTION',/22X,'(MB)',4X,'(METERS)',
26X,'(DEG K)',10X,'(M/SEC)',6X,'(DIRN. FROM)',//)
    ZE(1)=HT1-NDZ
    DO 290 I=1,N
    ZE(I)=ZE(1)+DZ
    TE(I)=TE(1)-CONST
    WRITE(6,300) PE(I),ZE(I),RHE(I),UE(I),THE(I)
    TE(I)=TE(I)+CONST
    CONTINUE
290 FORMAT(19X,F8.2,F11.2,F14.2,2PF15.3,OPF13.2,F17.0)
300 DO 320 I=1,N
310 ES=EXP(CX-AP/TE(I))/(TE(I)**8P)
    XS=(RATIO*ES)/(PE(I)-ES)
    XE(I)=XS*RHE(I)
    THE(I)=THE(1)*(PI/180.)
320 THE(I)=THE(I)*PI/180.)

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C ***** PART 2 MODEL COMPUTATIONS *****
C
C 330 NCR=1

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CS01410
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CS01500
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CS01540
CS01550
CS01560
CS01570
CS01580
CS01590
CS01600
CS01610
CS01620
CS01630
CS01640
CS01650
CS01660
CS01670
CS01680
CS01690
CS01700
CS01710
CS01720
CS01730
CS01740
CS01750
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CS01770
CS01780
CS01790
CS01800
CS01810
CS01820
CS01830
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CS01850
CS01860
CS01870
CS01880

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NBS=1
LWC=0
A1=0.2
AKF1=.001
AKF2=.0052
CRAD=1.0
TF=248.0
340  FORMAT(3I1,F4.0,6F10.0,F2.0)
350  SIZE(NUM)=CRAD
    AO(NUM)=A1
    NT=273.0-TF
    ZE(1)=IHT1
360  WRITE(6,360) AO(NUM),CRAD,AK1,AKF1,AK2,AKF2,TF
    FORMAT(1I10,F3.1,/,RUP,2X,RUP=,F6.3,/,KM,4X,K1=,F5.3,2X,
1  K1F=,F5.3,2X,K2=,F6.4,2X,K2F=,F6.4,4X,TF=,F6.1)
    RAD1=1000.0*CRAD
    XMUL=AO(NUM)/RAD1
    WRITE(6,370)
370  FORMAT(10,2X,HEIGHT,4X,HEIGHT,3X,PRESSURE,2X,VERTICAL,
13X,CLOUD,6X,TEMP,4X,MIXING,4X,CLOUD,5X,HYDRO,7X,Z,
28X,UPDRAFT,1)
    WRITE(6,380)
380  FORMAT(1,31X,VELOCITY,3X,TEMP,5X,EXCESS,4X,RATIO,5X,
1,WATER,5X,FACTOR,5X,RADIUS)
    WRITE(6,390)
390  FORMAT(1,1X,(METERS),3X,(FEET),5X,(MB),5X,(MPS),4X,
1,(DEG A),4X,(DEG C),2X,(GM/KG),3X,(GM/KG),3X,
2,(MM6/M3),4X,(METERS),/)
C ***** INITIALIZATION *****
C
C
410  IF(NBS) 420,420,410
    J=1
    NFRZ=0
    C NFRZ=KEY TO INDICATE WHEN FREEZING LEVEL IS REACHED
    TH(1)=THE(1)
    U(1)=UE(1)
    T(1)=TE(1)
    TC(1)=TE(1)
    C TC= TEMPERATURE OF CLOUD
    X(1)=XE(1)/RHE(1)
    QH(1)=0.0
    AQH(1)=0.0
    QCL(1)=0.0
    W(1)=0.5
    C W=0.5 IS ASSUMPTION OF .5 METERS PER SECOND VERTICAL VELOCITY

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CS01890
 CS01900
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 CS01980
 CS01990
 CS02000
 CS02010
 CS02020
 CS02030
 CS02040
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 CS02060
 CS02070
 CS02080
 CS02090
 CS02100
 CS02110
 CS02120
 CS02130
 CS02140
 CS02150
 CS02160
 CS02170
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 CS02200
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 CS02240
 CS02250
 CS02260
 CS02270
 CS02280
 CS02290
 CS02300
 CS02310
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 CS02330
 CS02340
 CS02350
 CS02360

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AW(1)=W(1)
IHT=IHT1
RO=0.0
RAD=RADI=CRAD
UPRAX=0.0
SDTMX=0.0
JDTMX=0
JSDTM=0
GO TO 440

C ***** INITIALIZATION AT PREVIOUS FREEZING LEVEL *****
C
C
420 J=ITOB
CTMAX=SDTMX
JDTMX=JSDTM
AREA(NUM)=AREA(NUM-1)
T(1)=TC(ITOB)
X(1)=AX(ITOB)
QCL(1)=AQC(ITOB)
QH(1)=AQH(ITOB)/1000.
W(1)=AW(ITOB)
RAD=UPRAD(ITOB)*RADI
RO=ROD
TH(1)=THTH
U(1)=UU
IHT=IHT1
IST=INDOC+1
DO 430 I=IST,ITOB,INDOC
IHT=IHT+NDZ*INDOC
IHTT=(IHT*328)/100
A=TC(I)-TE(I)
PRAD=UPRAD(I)*RADI
BQH=AQH(I)/1000.
C=(14000.*BQH)**1.136
ZFZC=200.*(C**1.6)
430 WRITE(6,790) IHT,IHTT,PE(I),AW(I),TC(I),A,AX(I),AQC(I),BQH,ZFZC,
1 PRAD
440 TVE(1)=T(1)*(1.0+.61*X(I))
Q(1)=QCL(1)+QH(1)
DEN(1)=PE(J)/(RD*TE(J))*1
INDIC=0
DA=0.0
C DA= HEAT ADDED DUE TO FREEZING OF LIQUID WATER
AP=6958.9262
RP=5.65567
CX=59.01383
XK4=15.39
  
```

```

AKA=AK2
AK=AK1
URASE=UE(1)
THBASE=THE(1)
IT=0
C   IT IS A KEY TO INDICATE WHEN ICE NUCLEATION LEVEL REACHED
450 J=J+1
451 IF(IT) 451,451,459
      XL=CON1*(T(1)-TNOT)+CON2
C   THIS CONVERTS XL FROM CALORIES/GRAM TO JOULES/KILOGRAM
459 IF(N-J) 810,460,460
460 XMT=AC(NUM)/RAD
      IHT=IHT+NDZ
      INDIC=INDIC+1
C
C ***** MOIST OR ICE ADIABATIC ASCENT *****
A=-G*DZ/CP
B=1.0+((X(1)*XL)/(RD*T(1)))
C=1.0+XL*X(1)/(CP*RV*T(1)**2)
C   SEE EQN. 190 PAGE 377 OF HANDBOOK OF METEOROLOGY
464 IF(IT) 464,464,465
      XL=CON1*(T(2)-TNOT)+CON2
      XL=XL*4.186E3
C
C ***** MIXING AT CONSTANT PRESSURE *****
465 ES=EXP(CX-AP/T(2))/(T(2)**BP)
      X(2)=(RATIO*ES)/(PE(J)-ES)
      A=(XMU1*DZ*XL/CP)*(X(2)-XE(J))
      B=(XMU*DZ)*(T(2)-TE(J))
      C=1.0+XL*X(2)/(CP*RV*T(2)**2)
C   A,B, AND C ARE NOW BEING REDEFINED (EQN.1 P.19 OF W.D. PAPER)
      T(2)=T(2)-(A+B)/C
      ES=EXP(CX-AP/T(2))/(T(2)**BP)
      X(2)=(RATIO*ES)/(PE(J)-ES)
C
C ***** CLOUD PHYSICS TERMS *****
      DENS=PE(J)/(RD*TE(J))*1.0E5
      AA=0.5000
      ADEN=AA/DENS
      IF(QCL(1)-ADEN) 470,470,480
470 A=0.0
      GO TO 490
C

```

CS02850
CS02860
CS02870
CS02880
CS02890
CS02900
CS02910
CS02920
CS02930
CS02940
CS02950
CS02960
CS02970
CS02980
CS02990
CS03000
CS03010
CS03020
CS03030
CS03040
CS03050
CS03060
CS03070
CS03080
CS03090
CS03100
CS03110
CS03120
CS03130
CS03140
CS03150
CS03160
CS03170
CS03180
CS03190
CS03200
CS03210
CS03220
CS03230
CS03240
CS03250
CS03260
CS03270
CS03280
CS03290
CS03300
CS03310
CS03320


```

750 IF(T(2)-TE(J)-T(1)+TE(J-1)) 750,770,770
760 HT=IHT-IHT1-NDZ
    DTMAX=T(1)-TE(J-1)
    AREA(NUM)=DTMAX/(T(1)-TE(1)+G/CP*HT)
    DENOM=T(1)-TE(1)+G/CP*HT
765 WRITE(6,765) AREA(NUM),DTMAX,DENOM
    FORMAT(2X,3E15.5,3X,'CARD CS03870',50(' '))
    JDTMX=J
77C IF(INDIC-INDOC) 800,780,780
780 INDIC=0
C C C
***** TABULAR OUTPUT OF PROFILES *****
D=T(2)-TE(J)
IHTFT=(IHT*328)/100
WRITE(6,790) IHT,IHTFT,PE(J),W(2),T(2),D,X(2),QCL(2),QH(2),ZFZC,
1RAD
79C FORMAT(18,I10, F11.4, F10.5,2F10.4,3P3F10.6,2X,CP,10.2,F10.2)
C C C
***** STORE PROFILES FOR GRAPHICAL OUTPUT *****
800 ITAB=J
    AW(ITAB)=W(2)
    TC(ITAB)=T(2)
    AQH(ITAB)=QH(2)*1000.0
    AX(ITAB)=X(2)
    AQC(ITAB)=QCL(2)
    UPRAD(ITAB)=RAD/RAD1
C C C
***** PREPARE FOR NEXT GRID STEP *****
Q(1)=Q(2)
TVE(1)=TVE(2)
T(1)=T(2)
X(1)=X(2)
QCL(1)=QCL(2)
QH(1)=QH(2)
DEN(1)=DEN(2)
W(1)=W(2)
GO TC 450
C C C
***** TOTAL PRECIPITATION *****
810 PA(NUM)=RO*39.37
C C C
***** DURATION OF PRECIPITATION *****

```

CS03800
CS03810
CS03820
CS03830
CS03840
CS03850
CS03860
CS03870
CS03880
CS03890
CS03900
CS03910
CS03920
CS03930
CS03940
CS03950
CS03960
CS03970
CS03980
CS03990
CS04000
CS04010
CS04020
CS04030
CS04040
CS04050
CS04060
CS04070
CS04080
CS04090
CS04100
CS04110
CS04120
CS04130
CS04140
CS04150
CS04160
CS04170
CS04180
CS04190
CS04200
CS04210
CS04220
CS04230
CS04240
CS04250
CS04260
CS04270


```

A=IHT-IHT1
IF(QH(2).GE..1E-5) GO TO 815
QH(2)=.1E-5
WRITE(6,814)
814 FORMAT(2X,'CARD CS04320 ',50(' '))
815 DUR(NUM)=(A/(XK4*QH(2)**.125))/60.0
***** CLOUD TOP HEIGHT *****
ITOP(NUM)=IHT
***** FREEZING TEMPERATURE (DEG. C) *****
NTF(NUM)=-NT
***** CLOUD TOP TEMPERATURE (DEG. C) *****
TEMPT(NUM)=T(2)-CONST
WRITE(6,820) RA(NUM), DUR(NUM), ITOP(NUM), AREA(NUM)
820 FORMAT(1X,'TOTAL RAIN=',F10.4,' INCHES PER CLOUD',5X,'RAIN LASTS',
1F10.2,' MINUTES',5X,'CLOUD TOP=',F15,' METERS',3X,'UPDRAFT AREA=',
22PF6.1,' %',//)
***** IF TOP DOES NOT REACH WARMEST FREEZING LEVEL, PROCEED TO *****
***** NEXT BOUNDARY CONDITION CARD *****
830 NCR=1
NBS=0
LWC=0
AI=0.2
AK1=.001
AKF1=.001
AK2=.0052
CRAD=1.0
TF=248.0
IF(NBS) 840,840,870
840 IF(NFRZ) 870,870,850
850 WRITE(6,860) AO(NUM), CRAD,AK1,AKF1,AK2,AKF2,TF
860 FORMAT(1X,'MU=',F3.1,' /RUP=',F5.3,' KM',4X,'K1=',F5.3,2X,
1,'K1F=',F5.3,2X,'K2=',F6.4,4X,'TF=',F6.1,2X,
2,'TOP TEMP. ABOVE TF,')
GO TO 900
***** GRAPHICAL OUTPUT OF SELECTED PROFILES *****
870 CALL GRPHCL(ITOP(NUM),NUM,NDZ,INDOC,ITAB,TC,TE,AW,AQH,UPRAD,AO(NUM
1))

```

CS04760
CS04770
CS04780

900 WRITE(6,820) RA(NUM),DUR(NUM),ITOP(NUM),AREA(NUM)
RETURN
END

```

C
C
C
SUBROUTINE GRPHCL(ITOP,NUM,NDZ,INDOC,ITAB,TC,TE,AW,AQH,UPRAD,A1)
SUBROUTINE TO GRAPH SELECTED PROFILES. IN THIS CASE PROFILES OF
VERTICAL VELOCITY, TEMPERATURE EXCESS, HYDROMETEOR WATER, AND
UPDRAFT RADIUS ARE GRAPHED.
INTEGER*2 CH,CHAR
DIMENSION TC(200),AW(200),AQH(200),UPRAD(200),VAR(10),
1STRT(10),GI(10),CH(10),Q(10),R(5),I(10),
DATA CH/1,W,2,Q,3,R,5*,I/
DATA STRT/-5,OE-1,-1,OE-1,-1,OE-1,-3.3333E-2,5*0.0E0/
DATA GI/5,OE0,2,OE0,1,OE0,2,OE0,3.3333E-1,5*0.0E0/
WRITE(6,5) ITAB,NUM,NDZ,A1
5 FORMAT(1,1,10X,'GRA P H I C A L D I S P L A Y NUMBER OF LEVEL
15=',I5,'//,11X,'BOUNDARY CONDITION NUMBER',I4,5X,'VERTICAL GRID DI
2STANCE=',I4,1X,'METERS',5X,'MU=',F5.2,'/RADIUS',//)
IHT=ITOP+NDZ*INDOC-NDZ
CALL SCALE(5,STRT,GI)
DO 10 I=1,ITAB,INDOC
IHT=IHT-NDZ*INDOC
J=ITAB-I+1
VAR(1)=AW(J)-TE(J)
VAR(2)=TC(J)-AQH(J)
VAR(3)=AQH(J)
VAR(4)=0.0
VAR(5)=UPRAD(J)
CALL GRAPH(5,CHAR,VAR,CH)
10 WRITE(6,20) IHT,CHAR
20 FORMAT(1HS,16,100A1)
30 WRITE(6,30) 7X,11('*,9X))
30 FORMAT(1H,7X,102('*,))
40 WRITE(6,40) 6X,102('*,))
40 FORMAT(1H,6X,102('*,))
50 WRITE(6,50) 7X,I=5,50,5)
50 FORMAT(1H,7X,I=5,50,5)
60 WRITE(6,60) 7X,I=2,10,2)
60 FORMAT(1H,7X,I=2,10,2)
1, T EXCESS(C))
70 WRITE(6,70) (I,I=1,10)
70 FORMAT(1H,7X,I=1,10)
80 WRITE(6,80) (I,I=1,3)
80 FORMAT(1H,7X,I=1,3)
80 FORMAT(8X,'0',3130,10X,' UPDRAFT RADIUS(KM).)
RETURN

```

GL000C10
GL000C20
GL000C30
GL000C40
GL000C50
GL000C60
GL000C70
GL000C80
GL000C90
GL000100
GL000110
GL000120
GL000130
GL000140
GL000150
GL000160
GL000170
GL000180
GL000190
GL000200
GL000210
GL000220
GL000230
GL000240
GL000250
GL000260
GL000270
GL000280
GL000290
GL000300
GL000310
GL000320
GL000330
GL000340
GL000350
GL000360
GL000370
GL000380
GL000390
GL000400
GL000410

FND

GL00420

```

SUBROUTINE GRAPH(N,CHAR,YC,CH)
  INTEGER*2 CH(10),CHAR(100),Z//',/'
  DIMENSION YC(10),GI(10),STRT(10),GIMOD(10)
  DO 10 I=1,100
    9 CHAR(I)=Z
    GO TO 30
  ENTRY SCALE(N,STRT,GI)
  DO 20 I=1,N
    20 GIMOD(I)=10.0/GI(I)
  RETURN
  30 DO 50 I=1,N
    DI=(YC(I)-STRT(I))*GIMOD(I)+1.0
    IS=DI
    IF((DI-IS).GE.0.5)IS=IS+1
    IF(IS.LE.0) GO TO 45
    40 IF(IS.LE.100) GO TO 42
    IS=IS-100
    GO TO 40
    42 CHAR(IS) = CH(I)
    GO TO 50
  45 WRITE(6,46) IS,DI,N,I,YC(I),STRT(I),GIMOD(I)
  46 FORMAT(2X,'IS=',I8,2X,'DI=',E9.1,2X,'N=',I3,2X,'I=',I3,2X,'YC=',
    1E9.1,2X,'STRT=',E9.1,2X,'GIMOD=',E9.1)
  50 CONTINUE
  RETURN
END

```

```

SUBROUTINE METMAP(Y,N,M,T,RND,AZ,BZ,AMIN,IJT,ICON)
  THIS SUBROUTINE OBJECTIVELY CONTOURS THE FIELD BROUGHT IN AS Y
  REAL*4 IH,KG,IJTJZ
  DIMENSION A(140),R(140),C(140),D(140),IH(20),Y(N,M),TP(5),XMT(5),
    1PT(5),KG(10),T(24)

  DATA DUE/4H /,EPL/4H+ /,EMI/4H- /,IH/1H0,1H /,1H1,1H /,1H2,
    11H /,1H3,1H /,1H4,1H /,1H5,1H /,1H6,1H /,1H7,1H /,1H8,1H /,1H9,1H /,KG/
    21H0,1H1,1H2,1H3,1H4,1H5,1H6,1H7,1H8,1H9/,BLK/4H /

  YMIN=Y(1,1)
  YMAX=Y(1,1)

```

```

      DO 20 I=1,M
      DO 10 J=1,N
      YMIN=AMIN1(YMIN,Y(J,I))
      YMAX=AMAX1(YMAX,Y(J,I))
10    CONTINUE
120   DELY=YMAX-YMIN
      IF(BND) 25,25,30
      RND=DELY/15.0
      IF (AMIN-YMIN) 31,31,32
      IF (IJT) 33,32,33
32   PD=YMIN/BND
      PF=ABS(PD-INT(PD))
      IF (YMIN) 21,1
      AMIN=YMIN-PF*BND
      GO TO 33
      2   AMIN=YMIN-(1.0-PF)*RND
33   AHLD=AZ
      IF(AZ) 55,35,55
35   SM=AMAX1(ABS(YMIN),ABS(YMAX))
      NS=0
40   NS=NS+1
      SM=10.0*SM
      IF(SM-1.0)40,50,45
45   NS=NS+1
      SM=SM/10.0
      IF(SM-1.0)50,50,45
50   AHLD=10.0*NS
55   HAND=BND/2.0
      PRINT 70
      PRINT 61
      FORMAT(5X,24A4,/)
      PRINT 57,AHLD,87
57   FORMAT(1H0,65H THE FOLLOWING TRANSFORMATION WAS PERFORMED ON THE IN
      IFUT MATRIX /5X,1H(E12.5,8H*(I,J)+,E12.5,1H)/2X,73HAND THREE
      2 DIGITS TO THE RIGHT OF THE DECIMAL POINT ARE PRINTED IN THE MAP )
      PRINT 54,YMAX,YMIN
54   FORMAT(/4X,5HYMAX=,E15.7,5X,5HYMIN=,E15.7)
      IF (ICON)5,58,5
5   PRINT 11,RND
11   FORMAT(2X,17H THE BAND WIDTH IS,E12.5,6H UNITS /4X,14H CONTOUR LEVE
      I=0
      YTOP=AMIN
      IF(ABS(YMIN-YMAX)-100.0*BND)53,53,58
53   Y8=YTOP
      YTOP=YTOP+RND

```

```

MP00140
MP00150
MP00160
MP00170
MP00180
MP00190
MP00200
MP00210
MP00220
MP00230
MP00240
MP00250
MP00260
MP00270
MP00280
MP00290
MP00300
MP00310
MP00320
MP00330
MP00340
MP00350
MP00360
MP00370
MP00380
MP00390
MP00400
MP00410
MP00420
MP00430
MP00440
MP00450
MP00460
MP00470
MP00480
MP00490
MP00500
MP00510
MP00520
MP00530
MP00540
MP00550
MP00560
MP00570
MP00580
MP00590
MP00600
MP00610

```

```

I=I+1
J=MOD(I,20)
ITJZ=IH(J)
IF(YB-YMAX)59,58,58
50 PRINT 61,Y8,YTOP,ITJZ
61 FORMAT(/4X,E10.3,4H TO ,E10.3,2H =,1X,A1)
58 NCCP=0
60 NCP=0
60 PRINT 70
70 FORMAT(IH1)
PRINT 6,T
NLINE=0
NCCP=NCP+1
NCP=NCP+25
73 IF(NCP-M)80,80,75
75 NCP=M
80 J=0
NLINE=NLINE+1
LLINE=N-NLINE+1
C SET UP HEADING
85 J=4
90 DO 100 I=1,135
A(I)=BLK
P(I)=BLK
C(I)=BLK
D(I)=BLK
100 CONTINUE
110 DO 160 L=NCCP,NCP
J=J+3
KI=L
IF(KI-100) 130,120,120
120 LL=KI/100
A(J)=KG(LL+1)
KI=KI-100*LL
GO TO 135
130 A(J)=KG(1)
135 J=J+1
IF(KI-10) 150,140,140
140 LL=KI/10
A(J)=KG(LL+1)
KI=KI-10*LL
GO TO 155
150 A(J)=KG(1)
155 J=J+1
A(J)=KG(KI+1)
160 CONTINUE

```

```

MP00620
MP00630
MP00640
MP00650
MP00660
MP00670
MP00680
MP00690
MP00700
MP00710
MP00720
MP00730
MP00740
MP00750
MP00760
MP00770
MP00780
MP00790
MP00800
MP00810
MP00820
MP00830
MP00840
MP00850
MP00860
MP00870
MP00880
MP00890
MP00900
MP00910
MP00920
MP00930
MP00940
MP00950
MP00960
MP00970
MP00980
MP00990
MP01000
MP01010
MP01020
MP01030
MP01040
MP01050
MP01060
MP01070
MP01080
MP01090

```

```

C SETUP FIRST ROW OF ARRAY
170 GO TO 260
    NLINE=NLINE+1
    LLINE=N-NLINE+1
    IF(NLINE-N) 180,180,380
180 DO 190 I=1,135
    A(I)=RLK
    R(I)=BLK
    C(I)=BLK
    CONTINUE
190 IF (ICON)195,260,195
195 NCV=NCCP-1
    J=1
    IF(NCY)200,200,210
200 J=5
210 NCV=NCY+1
    IF(NCY-NCP) 220,220,260
    IF(NCY-M) 230,260,260
220 NLINE = NLINE - 1
230 YD1=Y(NLINE,NCY)-Y(NLINE+1,NCY)
    YD2=Y(NLINE,NCY+1)-Y(NLINE+1,NCY+1)
    TP(1)=Y(NLINE,NCY)-0.25*YD1
    XMT(1)=Y(NLINE,NCY)-0.5*YD1
    RT(1)=Y(NLINE,NCY)-0.75*YD1
    TP(5)=Y(NLINE,NCY+1)-0.25*YD2
    XMT(5)=Y(NLINE,NCY+1)-0.5*YD2
    RT(5)=Y(NLINE,NCY+1)-0.75*YD2
    NLINE = NLINE + 1
    C1=0.25*(TP(5)-TP(1))
    C2=0.25*(XMT(5)-XMT(1))
    C3=0.25*(RT(5)-RT(1))
    DO 240 I=2,4
    TP(I)=TP(I-1)+D1
    XMT(I)=XMT(I-1)+D2
    RT(I)=RT(I-1)+D3
240 CONTINUE
    J=J+1
    IF(J=1) 250
    I1=MOD(IFIX((TP(I)-AMIN)/BND),20)+1
    I2=MOD(IFIX((XMT(I)-AMIN)/BND),20)+1
    I3=MOD(IFIX((RT(I)-AMIN)/AND),20)+1
    A(J)=IH(I1)
    R(J)=IH(I2)
    C(J)=IH(I3)
    CONTINUE
250 GO TO 210
260 NCV=NCCP-1

```

MP01100
 MP01110
 MP01120
 MP01130
 MP01140
 MP01150
 MP01160
 MP01170
 MP01180
 MP01190
 MP01200
 MP01210
 MP01220
 MP01230
 MP01240
 MP01250
 MP01260
 MP01270
 MP01280
 MP01290
 MP01300
 MP01310
 MP01320
 MP01330
 MP01340
 MP01350
 MP01360
 MP01370
 MP01380
 MP01390
 MP01400
 MP01410
 MP01420
 MP01430
 MP01440
 MP01450
 MP01460
 MP01470
 MP01480
 MP01490
 MP01500
 MP01510
 MP01520
 MP01530
 MP01540
 MP01550
 MP01560
 MP01570

```

265 J=0
265 IF(NCY) 265,265,270
270 GO TO 330
270 NCY=NCY+1
280 IF(NCY-NCP) 280,280,310
280 J=J+2
285 THLD=AHLD*Y(NLINE,NCY)+87
285 IF(THLD) 285,290,290
285 C(J)=EMI
290 GO TO 295
290 C(J)=EPL
295 NUM=INT(ABS(THLD-INT(THLD)))*1000.0+0.5)
295 NDS=100
300 KK=1,3
300 J=J+1
300 KI=NUM/NDS
300 C(J)=KG(KI+1)
300 NUM=NUM-KI*NDS
300 NDS=NDS/10
300 CONTINUE
300 GO TO 270
310 IF(NCP-M) 360,320,320
320 IF(J-127)330,330,360
330 J=J+3
335 KI=NLINE
335 IF(KI-100) 340,335,335
335 LL=KI/100
335 C(J)=KG(LL+1)
335 KI=KI-100*LL
340 GO TO 343
340 C(J)=KG(1)
343 J=J+1
343 IF(KI-10) 350,345,345
345 LL=KI/10
345 C(J)=KG(LL+1)
345 KI=KI-10*LL
350 GO TO 355
350 C(J)=KG(1)
355 J=J+1
355 C(J)=KG(KI+1)
355 IF(NCY-1) 270,270,360
360 IF(NLINE-1)362,362,368
362 PRINT 370,(A(I),I=1,132),(B(IP1),IP1=1,132),(D(IP2),IP2=1,132)
368 GO TO 170
368 PRINT 370,(A(I),I=1,132),(B(IP1),IP1=1,132),(C(IP2),IP2=1,132),
370 1(D(IP3),IP3=1,132)
370 FORMAT(132A1)

```

```

MP01580
MP01590
MP01600
MP01610
MP01620
MP01630
MP01640
MP01650
MP01660
MP01670
MP01680
MP01690
MP01700
MP01710
MP01720
MP01730
MP01740
MP01750
MP01760
MP01770
MP01780
MP01790
MP01800
MP01810
MP01820
MP01830
MP01840
MP01850
MP01860
MP01870
MP01880
MP01890
MP01900
MP01910
MP01920
MP01930
MP01940
MP01950
MP01960
MP01970
MP01980
MP01990
MP02000
MP02010
MP02020
MP02030
MP02040
MP02050

```

```

380      CN TN 170
      A(I)=RLK I=1,135
      R(I)=RLK
      C(I)=RLK
      D(I)=RLK
390      CONTINUE
      J=0
395      IF(NCCP-1) 395,395,400
400      J=4
      CN 430 L=NCCP,NCP
      J=J+3
      KI=L
      IF(KI-100) 410,405,405
405      LL=KI/100
      C(J)=KG(LL+1)
      KI=KI-100*LL
      CN TN 412
410      C(J)=KG(1)
412      J=J+1
      IF(KI-10) 420,415,415
415      LL=KI/10
      C(J)=KG(LL+1)
      KI=KI-10*LL
      CN TN 422
420      C(J)=KG(1)
422      J=J+1
      IF(KI-1) 430,415,415
430      CONTINUE
      PRINT 370, (B(IP1), IP1=1,132), (C(IP2), IP2=1,132)
      IF(NCP-M)60,500,500
500      PRTURN
      END

```

```

MP02060
MP02070
MP02080
MP02090
MP02100
MP02110
MP02120
MP02130
MP02140
MP02150
MP02160
MP02170
MP02180
MP02190
MP02200
MP02210
MP02220
MP02230
MP02240
MP02250
MP02260
MP02270
MP02280
MP02290
MP02300
MP02310
MP02320
MP02330
MP02340
MP02350
MP02360
MP02370
MP02380

```



```

C      THIS IS A PROGRAM TO DO AN OBJECTIVE ANALYSIS FOR NSSL NETWORK.
C      X-AXIS POINTS EAST AND Y-AXIS POINTS NORTH. POINT 1,1 IS LOWER
C      LEFT CORNER OF THE GRID
C
C      DIMENSION Z(24,24), ZDUM(24,24), TITLE(24), XTRA(24,24)
C      Z = FIELD TO BE ANALYZED
C      ZDUM = FIELD TO USE IN METMAP
C      TITLE = TITLE TO USE IN METMAP
C      XTRA = FIELD TO STORE ORIGINAL Z-FIELD IF SEVERAL VALUES OF 'C'
C      ARE TO BE TESTED
C      DO 975 NSETS=1,2
C      TWO CARDS NEED TO BE CHANGED EACH RUN. THEY ARE THE PRECEEDING
C      CARD THAT DETERMINES THE NUMBER OF SETS OF DATA THAT ARE TO BE
C      ANALYZED, AND THE WRITE(7, ) STATEMENT THAT DETERMINES WHAT
C      FORMAT IS TO BE USED ON THE PUNCHED CARDS. FOUR DATA CARDS ARE
C      REQUIRED FOR EACH SET OF DATA TO BE ANALYZED. THE NINE ZGES
C      VALUES ARE FOR THE 'REGIONS OF INFLUENCE'. BND AND SCALE ARE
C      FOR CALLING ARGUMENTS FOR SUBROUTINE METMAP. C = THE ARBITRARY
C      CONSTANT IN THE LAPLACIAN PART OF THE ANALYSIS SCHEME.
C      NI=24
C      NJ=24
C      NI = NUMBER OF GRID POINTS ALONG THE X-AXIS
C      NJ = NUMBER OF GRID POINTS ALONG THE Y-AXIS
C      NI=NI-1
C      NJ=NJ-1
C      READ(5,10) TITLE
C      FORMAT(20A4)
C      READ(5,20) ZGES1,ZGES2,ZGES3,ZGES4,ZGES5,ZGES6,ZGES7,ZGES8,ZGES9,
C      BND,C,SCALE
C      FORMAT(6F10.4)
C      ZGES=(ZGES1+ZGES2+ZGES3+ZGES4+ZGES5+ZGES6+ZGES7+ZGES8+ZGES9)/9.C
C      DO 22 I=1,14
C      DO 22 J=1,15
C      Z(I,J) = ZGES1
C      XTRA(I,J)=Z(I,J)
C      22 CONTINUE
C      DO 23 I=15,24
C      DO 23 J=1,15
C      Z(I,J) = ZGES2
C      XTRA(I,J)=Z(I,J)
C      23 CONTINUE
C      DO 24 I=1,6
C      DO 24 J=6,11
C      Z(I,J) = ZGES3
C      XTRA(I,J)=Z(I,J)
C      24 CONTINUE
C      DO 25 I=7,17

```



```

DO 25 J=6,11
  Z(I,J) = ZGES4
  XTRA(I,J)=Z(I,J)
25 CONTINUE
DO 26 I=18,24
  DO 26 J=6,11
    Z(I,J) = ZGES5
  XTRA(I,J)=Z(I,J)
26 CONTINUE
DO 27 I=1,10
  DO 27 J=12,20
    Z(I,J) = ZGES6
  XTRA(I,J)=Z(I,J)
27 CONTINUE
DO 28 I=11,17
  DO 28 J=12,20
    Z(I,J) = ZGES7
  XTRA(I,J)=Z(I,J)
28 CONTINUE
DO 29 I=18,24
  DO 29 J=12,24
    Z(I,J) = ZGES8
  XTRA(I,J)=Z(I,J)
29 CONTINUE
DO 30 I=1,17
  DO 30 J=21,24
    Z(I,J) = ZGES9
  XTRA(I,J)=Z(I,J)
30 CONTINUE
C NOW PERFORM THE ANALYSIS SCHEME
70 BOUND=Z(10,1)
DO 150 LPASS=1,11 GO TO 124
DO 110 LPLAC=1,10
DO 100 I=2,NII
DO 100 J=2,NJI
C THE FOLLOWING 'IF' STATEMENTS ARE NECESSARY SO AS NOT TO ALTER THE
C VALUES OF THE KNOWN GRID POINTS
C
  IF(I.EQ.10.AND..J.EQ. 1) GO TO 100
  IF(I.EQ.19.AND..J.EQ. 3) GO TO 100
  IF(I.EQ.11.AND..J.EQ. 9) GO TO 100
  IF(I.EQ. 2.AND..J.EQ. 9) GO TO 100
  IF(I.EQ.23.AND..J.EQ. 9) GO TO 100
  IF(I.EQ.15.AND..J.EQ.14) GO TO 100
  IF(I.EQ. 5.AND..J.EQ.17) GO TO 100
  IF(I.EQ.21.AND..J.EQ.18) GO TO 100

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0A00490
0A00500
0A00510
0A00520
0A00530
0A00540
0A00550
0A00560
0A00570
0A00580
0A00590
0A00600
0A00610
0A00620
0A00630
0A00640
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0A00660
0A00670
0A00680
0A00690
0A00700
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0A00900
0A00910
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0A00940
0A00950
0A00960

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      IF(I.EQ.11.AND.J.EQ.23) GO TO 100
      PL=Z(I+1,J)+Z(I-1,J)+Z(I,J+1)+Z(I,J-1)-4.0*Z(I,J)
      Z(I,J)=Z(I,J)+C*PL
100 CONTINUE
110 CONTINUE
C
C NOW DO THE EIGHT POINT AVERAGE
C
124 DO 125 J=2,NJ1
    DO 125 I=2,NI1
      IF(I.EQ.10.AND.J.EQ.1) GO TO 125
      IF(I.EQ.19.AND.J.EQ.3) GO TO 125
      IF(I.EQ.11.AND.J.EQ.9) GO TO 125
      IF(I.EQ.2.AND.J.EQ.9) GO TO 125
      IF(I.EQ.23.AND.J.EQ.9) GO TO 125
      IF(I.EQ.15.AND.J.EQ.14) GO TO 125
      IF(I.EQ.5.AND.J.EQ.17) GO TO 125
      IF(I.EQ.21.AND.J.EQ.18) GO TO 125
      IF(I.EQ.11.AND.J.EQ.23) GO TO 125
      Z(I,J)=(Z(I+1,J)+Z(I-1,J)+Z(I,J+1)+Z(I,J-1)+
      125 CONTINUE
      1Z(I-1,J+1)+Z(I-1,J-1)+Z(I+1,J+1)+Z(I+1,J-1))/8.0
C
C NOW SET BOUNDARY CONDITIONS TO SAME VALUE AS NEAREST ROW OR COLUMN
C
    DO 135 I=1,NI
    DO 135 J=1,NJ
      Z(1,J)=Z(2,J)
      Z(I,1)=Z(I,2)
      Z(NI,J)=Z(NI-1,J)
      Z(I,NJ)=Z(I,NJ-1)
135 CONTINUE
      Z(10,1)=8000
150 CONTINUE
C
C NOW HAVE ANALYZED FIELD
C NOW SET UP FIELD TO USE IN METMAP
C
155 DO 250 I=1,NI
    DO 250 J=1,NJ
      ZDUM(I,J)=Z(I,NJ-J+1)
250 CONTINUE
      CALL METMAP(7DUM,NI,NJ,TITLE,BND,SCALE,0.0,ZGES,0,1)
310 FORMAT(12F6.1)
311 FORMAT(12F6.0)
      C=C+.01
      IF(C.GT.1.0) GO TO 900
    DO 350 I=1,NI

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OA00970
 OA00980
 OA00990
 OA01000
 OA01010
 OA01020
 OA01030
 OA01040
 OA01050
 OA01060
 OA01070
 OA01080
 OA01090
 OA01100
 OA01110
 OA01120
 OA01130
 OA01140
 OA01150
 OA01160
 OA01170
 OA01180
 OA01190
 OA01200
 OA01210
 OA01220
 OA01230
 OA01240
 OA01250
 OA01260
 OA01270
 OA01280
 OA01290
 OA01300
 OA01310
 OA01320
 OA01330
 OA01340
 OA01350
 OA01360
 OA01370
 OA01380
 OA01390
 OA01400
 OA01410
 OA01420
 OA01430
 OA01440

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DO 350 J=1,NJ
  Z(I,J)=XTRA(I,J)
  350 CONTINUE
  GO TO 70
500 WRITE(7,311) ((Z(I,J),I=1,NI),J=1,NJ)
575 CONTINUE
599 STOP
END

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0A01450
0A01460
0A01470
0A01480
0A01490
0A01500
0A01510
0A01520

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<p>A numerical model which utilizes the isobaric vorticity equation is developed and applied to cumulus-scale data. The model, together with a modified version of the cumulus convection model of Weinstein and Davis, is applied to data obtained from the National Severe Storms Laboratory in Norman, Oklahoma. The calculations yield real time predictions for height, temperature and relative humidity at seven pressure levels, which are then used as input to the cumulus convection model to obtain vertical profiles of various parameters at specified grid points.</p> <p>Some results of the calculations are presented along with suggestions for further testing and improvement. The results indicate that further modifications to the approach used are necessary in order to provide more accurate forecasts. Values of the individual terms in the vorticity equation are presented as computed from the observed mesoscale data.</p>			

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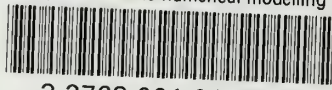
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Cumulus-scale motions



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An approach to the numerical modelling o



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